Programming incommensurately

Lukáš Palatinus and Václav Petříček Institute of Physics ASCR Praha, Czech Republic

A revolutionary proposal

- Rhombohedral space groups are notoriously difficult to deal with.
- Their symmetry properties are different from other crystals, and require special handling
- They account only for about 1% of all structures (according to CSD)

Let us just discard rhombohedral space groups! Programs will be easier to write and less buggy. And that 1% of structures? We have already solved about a milion structures, so who cares?

A revolutionary proposal

- Incommensurate structures are notoriously difficult to deal with.
- Their symmetry properties are different from other crystals, and require special handling
- They account only for about 1% of all structures (wild guess here)

Let us just discard incommensurate structures! Programs will be easier to write and less buggy. And that 1% of structures? We have already solved about a milion structures, so who cares?

Superspace at a glance











- When dealing with an incommensurately modulated structure we face four main challenges:
- Handling the diffraction pattern
- Handling the symmetry
- Computing the structure factor
- Computing structural properties (distances, angles...)

Common feature for all issues related to modulated structures: the use of superspace and work in more than 3 dimensions, but keeping in mind the relationship to 3D.

Diffraction pattern

3D

Properties:

- indices hkl
- diffraction vector
 h=ha*+kb*+lc*
- intensity I(h)
- structure factor F(h)
- d-spacing d=1/|h|
- d_{min} uniquely determines a set of reflections within certain resolution sphere

4D

Properties:

- indices hklm
- diffraction vector
 h=h*a**+k*b**+l*c**+m*q*
- intensity I(h)
- structure factor F(h)
- d-spacing d=1/|h|
- d_{min} and m_{max} determine a set of reflections within certain resolution sphere

Symmetry

3D

3D symmetry operators: $S = (\mathbf{R}, \mathbf{t})$ space group $G = \{S_i\}$ transformation of a vector: $\mathbf{r}' = \mathbf{Rr} + \mathbf{t} \cdot \mathbf{r}$ transformation of an atom $\mathbf{r}' = \mathbf{Rr} + \mathbf{t} \cdot \mathbf{r}$

4D

4D symmetry operators: $S = (\mathbf{R}, \mathbf{t})$ space group $G = \{S_i\}$ transformation of a vector: $\mathbf{r}' = \mathbf{Rr} + \mathbf{t} \cdot \mathbf{r}$ transformation of an atom $\mathbf{r} = \mathbf{r}_0 + \mathbf{u}(x_4)$ $\mathbf{r}_0' = \mathbf{R}_{\mathrm{E}} \mathbf{r}_0 + \mathbf{t}_{\mathrm{E}} \cdot \mathbf{r}_0$ $\mathbf{u}'(x_4) = \mathbf{R}_3 \mathbf{u}(R_{44}(x_4 - t_4))$

General symmetry element:
$$S = \begin{pmatrix} \cong \mathbf{R}_E & \mathbf{0} \cong \cong \mathbf{t}_E \cong \\ \Xi \mathbf{R}_M & \mathbf{R}_I \cong \oplus \mathbf{t}_I \cong \\ \end{pmatrix}$$

General symmetry element:
$$S = \begin{pmatrix} \cong \mathbf{R}_E & \mathbf{0} \\ \cong \mathbf{R}_M & \mathbf{R}_I \\ \cong \mathbf{R}_M & \mathbf{R}_I \\ \cong \mathbf{C} \\ \cong \mathbf{R}_I \\ \cong \mathbf{C} \\$$

1. The 3D symmetry element as determined from the main reflections



- 1. The 3D symmetry element as determined from the main reflections
- 2. Internal space cannot be mixed up with external space

General symmetry element:
$$S = \begin{pmatrix} \cong \mathbb{R}_E & \mathbb{0} \\ \cong \mathbb{R}_M & \mathbb{R}_I \\ \cong \mathbb{R}_I & \cong \mathbb{R}_I \\ \cong \mathbb{R}_I$$

- 1. The 3D symmetry element as determined from the main reflections
- 2. Internal space cannot be mixed up with external space
- 3. Form the metric properties \rightarrow **qR**_{*E*} **R**_{*I*}**q** = **R**_{*M*}

General symmetry element:
$$S = \begin{pmatrix} \exists \mathbf{R}_E & \mathbf{0} \\ \exists \mathbf{R}_M & \mathbf{R}_I \\ \exists \mathbf{R}_I & \exists \mathbf{r}_I \\ \exists \mathbf{r$$

- 1. The 3D symmetry element as determined from the main reflections
- 2. Internal space cannot be mixed up with external space
- 3. Form the metric properties \rightarrow **qR**_{*E*} **R**_{*I*}**q** = **R**_{*M*}

The only new information is the intrinsic part of the 4th component of the translation vector which, analogically to 3D symmetry, affects systematic absences of reflections. It expresses, how the modulation wave is shifted in the internal space. If \mathbf{t}_{i} is intrinsic, it can be determined from systematic absences.

Systematic absences

3D

- Determine the invariant reflections of the rotation matrix from the relationship h_{inv}=h_{inv}R
- 2) The invariant reflections match the relationship

h.t=*n*

Example: c-glide perp. to **b**:

 $x_1, x_2, x_3 \rightarrow x_1, -x_2, \frac{1}{2} + x_3$ $\mathbf{h}_{inv} = (h, 0, l); \mathbf{h} \cdot \mathbf{t} = l/2 = n:$

l=2n

4D

1) Determine the invariant reflections of the rotation matrix from the relationship

 $\mathbf{h}_{inv} = \mathbf{h}_{inv} \mathbf{R}$

2) The invariant reflections match the relationship

h.**t**=*n*

Example: c|*s*-glide perp. to **b**:

 $X_1, X_2, X_3, X_4 \rightarrow X_1, -X_2, \frac{1}{2} + X_3, \frac{1}{2} + X_4$

 $\mathbf{h}_{inv} = (h, 0, l, m); \mathbf{h} \cdot \mathbf{t} = l/2 + m/2 = n:$

l+m=2n

Systematic absences: *I+m=2n*

zero layer

first layer





Structure factor

3D

Basic formula:

$$F_{\mathbf{h}} = \int_{V} \rho(\mathbf{r}) \exp(2\pi i \mathbf{h} \cdot \mathbf{r}) \mathrm{d}V$$

Integrating out the atomic densities:

$$F_{\mathbf{h}} = \sum_{j} f_{j}(h) \exp(2\pi i \mathbf{h} \cdot \mathbf{r}_{j})$$

4D

Basic formula:

$$F_{\mathbf{h}} = \int_{V} \rho(\mathbf{r}) \exp(2\pi i \mathbf{h} \cdot \mathbf{r}) \mathrm{d}V$$

Integrating out the atomic densities:

$$F_{\mathbf{h}} = \sum_{j} f_{j}(h) \exp(2\pi i \mathbf{h} \cdot \mathbf{r}_{j}) \times \\ \times \int_{x_{4}} \exp(2\pi i (u(x_{4}) + mx_{4})) dx_{4}$$

Structure factor

$$F_{\mathbf{h}} = \sum_{j} f_{j}(h) \exp(2\pi i \mathbf{h} \cdot \mathbf{r}_{j}) \times \\ \times \int_{x_{4}} \exp(2\pi i (u(x_{4}) + mx_{4})) dx_{4}$$

For simple modulations (straight lines, harmonic modulation functions) analytical formulae exist.

Various computing methods have been devised for fast calculation of the structure factor, but nowadays the "simple" numerical integration (Gaussian quadrature) is used.

Poor man's solution: calculate superspace electron density on a regular grid, and get structure factors by FFT of the electron density

If you intend to make your program applicable to 3+n dimensions: 1)Do not fix the number of coordinates of your vectors to three. Make all vectors n-dimensional instead.

```
This brings two main problems:
```

```
a) constructs like this one are impossible:
    D0 h=hmin,hmax
    D0 k=kmin,kmax
    D0 l=lmin,lmax
        some clever math here
        ENDDO
ENDDO
ENDDO
```

If you intend to make your program applicable to 3+n dimensions: 1)Do not fix the number of coordinates of your vectors to three. Make all vectors n-dimensional instead.

This brings two main problems:

a) instead you need:

```
setup starting and ending coordinates
DO
CALL GiveMeNextIndices(hkl,IsLast)
some clever math here
IF (IsLast) EXIT
```

ENDDO

If you intend to make your program applicable to 3+n dimensions: 1)Do not fix the number of coordinates of your vectors to three. Make all vectors n-dimensional instead.

This brings two main problems:

b) multidimensional arrays (like electron density) must be stored in one-dimensional array. You have to write functions for transforming the index of the 1D array to the multidimensional indices and vice versa.

If you intend to make your program applicable to 3+n dimensions:

2) Avoid tables, especially symmetry tables. Use matrix algebra instead:

- it is more elegant
- does not suffer from the problem of non-standard settings
- tables for higher dimensions are very difficult to assemble
- algebraic approach is less prone to typing errors

If you intend to make your program applicable to 3+n dimensions: 2)Avoid tables, especially symmetry tables. Use matrix algebra instead:

Symmetry operations compatible with the lattice and centering:

	agreem	ent factor			
n(0,1,0):	1/2+x1	-x2	1/2+x3	0.116	
b(1,0,0):	-x1	1/2+x2	x 3	0.184	
2_1(0,0,1):	-x1	-x2	1/2+x3	0.237	
2_1(1,0,0):	1/2+x1	-x2	- x 3	33.107	
-1:	-x1	-x2	- x 3	33.207	
m(0,0,1):	x 1	x 2	-x3	75.603	
2_1(0,1,0):	-x1	1/2+x2	-x3	75.670	
m(0,1,0):	x 1	-x2	x 3	82.962	

Space group derived from the symmetry operations:

		HM symbol:	
		2c -2ab	Hall symbol:
	(3/4,3/4,0)	300223}250qY2	Fingerprint:
		ons:	Symmetry operat
х3	x 2	x1	1:
1/2+x3	-x2	-x1	2_1(0,0,1):
x 3	1/2+x2	1/2-x1	b(1,0,0):
1/2+x3	1/2-x2	1/2+x1	n(0,1,0):

If you intend to make your program aplicable to 3+n dimensions: 2)Avoid tables, especially symmetry tables. Use matrix algebra instead:

Symmetry operations compatible with the lattice and centering:

Symmetry operation				
1/2+x1	-x2	1/2+x3	0.116	
-x1	1/2+x2	ж3	0.184	
-x1	-x2	1/2+x3	0.237	
1/2+x1	-x2	-x3	33.107	
-x1	-x2	-x3	33.207	
x 1	x 2	-x3	75.603	
-x1	1/2+x2	-x3	75.670	
x1	-x2	ж3	82.962	
	Symme 1/2+x1 -x1 -x1 1/2+x1 -x1 x1 -x1 x1	Symmetry operation 1/2+x1 -x2 -x1 1/2+x2 -x1 -x2 1/2+x1 -x2 -x1 -x2 -x1 -x2 -x1 1/2+x2 -x1 1/2+x2 x1 x2 -x1 1/2+x2 x1 -x2	Symmetry operation 1/2+x1 -x2 1/2+x3 -x1 1/2+x2 x3 -x1 -x2 1/2+x3 1/2+x1 -x2 -x3 -x1 -x2 -x3 -x1 -x2 -x3 x1 x2 -x3 x1 x2 -x3 x1 1/2+x2 -x3 x1 -x2 x3 x1 -x2 x3	Symmetry operation agreement 1/2+x1 -x2 1/2+x3 0.116 -x1 1/2+x2 x3 0.184 -x1 -x2 1/2+x3 0.237 1/2+x1 -x2 -x3 33.107 -x1 -x2 -x3 33.207 x1 x2 -x3 75.603 -x1 1/2+x2 -x3 75.670 x1 -x2 x3 82.962

Space group derived from the symmetry operations:

HM symbol: Pbn21 Hall symbol: P 2c -2ab Fingerprint: 3300223}250qY2 (3/4,3/4,0) Symmetry operations: 1: x2 **x**3 x1 -x1 2 1(0,0,1):-x2 1/2 + x3 $1/2 + x^2$ хЗ b(1,0,0): 1/2-x1 n(0,1,0): 1/2+x1 $1/2 - x^2$ 1/2 + x3

Useful exercise

You may try to practice some of the superspace concepts by writing the following program:

Given a superspace group (as a list of one-line representations of **R** and **t**), produce a list of all systematically absent reflections within a given resolution sphere and given maximal satellite index.