

# XANES: MXAN and FPMS codes

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# Acknowledgement

- **Isabella Ascone**  
**Grazie mille!**
- **Joaquin García Ruiz**  
**¡Muchas gracias!**
- **Hiroyuki Oyanagi**  
**どうもありがとう御座います。**

**Many thanks!**

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Multiple scattering approach

**KKR** (band): Korringa (1947), **Kohn** and Rostoker (1954)

**SCF-SW** (Quantum Chemistry:MO): **Slater** and Johnson (1966),  
Smith and Johnson (1969)

For spectroscopy,

LEED: Pendry

XPD : Rennert, Natoli & Sebilliau

XAFS: Natoli (1980), Rehr (FEFF code), Benfatto, Brouder, Foulis,  
Ebert, Fujikawa, Siper, ...

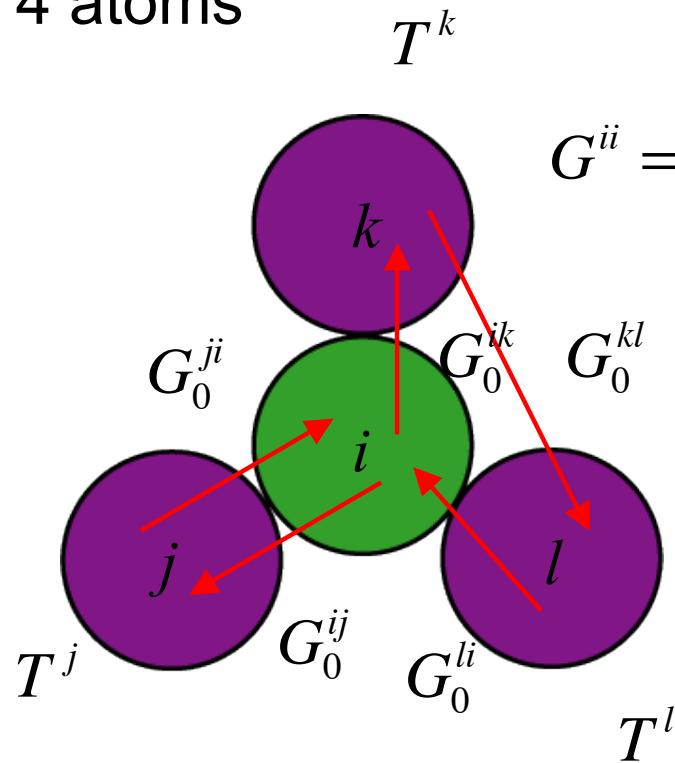
Also for ELNES, RXS

wide application of MS theory

K. Hatada, *IXAS Web Magazine*, 2nd issue, July (2010)

# Multiple scattering with Muffin-tin approximation : Green's function picture

4 atoms



Obtain  $G^{ii}$ :

$$G^{ii} = \underline{G_0^{ii}} + \underline{\sum_n G_0^{in} T^n G_0^{ni}} + \underline{\sum_{nm} G_0^{in} T^n G_0^{nm} T^m G_0^{mi}} + \dots$$

$G_0^{ii}$  : isolated

$\sum_n G_0^{in} T^n G_0^{ni}$  : single scattering

$\sum_{nm} G_0^{in} T^n G_0^{nm} T^m G_0^{mi}$

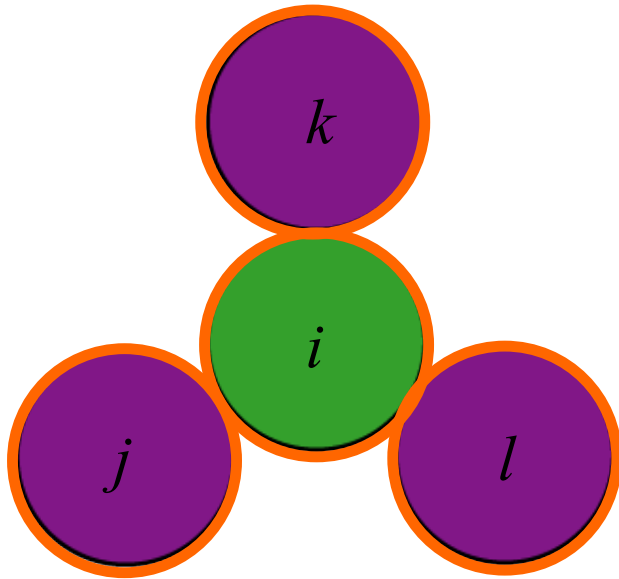
: double scattering

$$\begin{aligned} G^{ii} &= G_0^{ii} + \sum_{n \neq i} G_0^{in} T^n G_0^{ni} + \sum_{n \neq m \neq i} G_0^{in} T^n G_0^{nm} T^m G_0^{mi} + \dots \\ &= [G_0 (I - G_0 T)^{-1}]^{ii} \end{aligned}$$

Full Multiple Scattering

## Multiple scattering : wave function (WF) picture

WF is expanded by local solutions



Matches WF and its derivative on the surface of cells!

$$\text{Site } i: \Psi(\mathbf{r}_i; \mathbf{k}) = \sum_L B_L^i(\mathbf{k}) \Phi_L^i(\mathbf{r}_i)$$

$$\text{Site } j: \Psi(\mathbf{r}_j; \mathbf{k}) = \sum_L B_L^j(\mathbf{k}) \Phi_L^j(\mathbf{r}_j)$$

$$\text{Site } k: \Psi(\mathbf{r}_k; \mathbf{k}) = \sum_L B_L^k(\mathbf{k}) \Phi_L^k(\mathbf{r}_k)$$

$$\text{Site } l: \Psi(\mathbf{r}_l; \mathbf{k}) = \sum_L B_L^l(\mathbf{k}) \Phi_L^l(\mathbf{r}_l)$$

$B_L^i(\mathbf{k})$  : amplitude

$\Phi_L(\mathbf{r}_i)$  : local solution of SE

$\Psi(\mathbf{r}_i; \mathbf{k})$ : global solution of SE

$$\sum_{jL'} [T^{-1}(I - TG_0)]_{LL'}^{ij} B_{L'}^j = I_L^i(\mathbf{k})$$

This gives amplitudes on each sites. In other words matching will be done implicitly by solving this equation.

$I_L^i(\mathbf{k})$  :Boundary condition of the cluster for the cite i centered.



**MXAN** code: fitting structural and electronic properties by XANES spectra for molecule, protein and amorphous. (disordered system)

M. Benfatto and S. Della Longa, *J. Synchrotron Rad.*, **8**, 1087 (2001)

- XANES (X-ray Absorption Near Edge Structure) from edge to ~100 eV
- 3 dimensional information, sensitive to atomic species

Many applications have been done.

- S. Della Longa et al. Biophys. Jour. 85, 549 (2003)
- P. Frank et al. Inorganic Chemistry 44, 1922 (2005)
- C. Monesi et al. PRB 72, 174104 (2005)
- R. Sarangi et al. Inorganic Chemistry 44, 9652 (2005)
- P. D'Angelo et al. JACS 128, 1853 (2006)
- S. Marino et al. Biophys. Jour. 93, 2781 (2007)
- ....
- M. Bortolus et al. JACS 132, 18057 (2010)

Two ways to calculate the scattering path operator

by series:

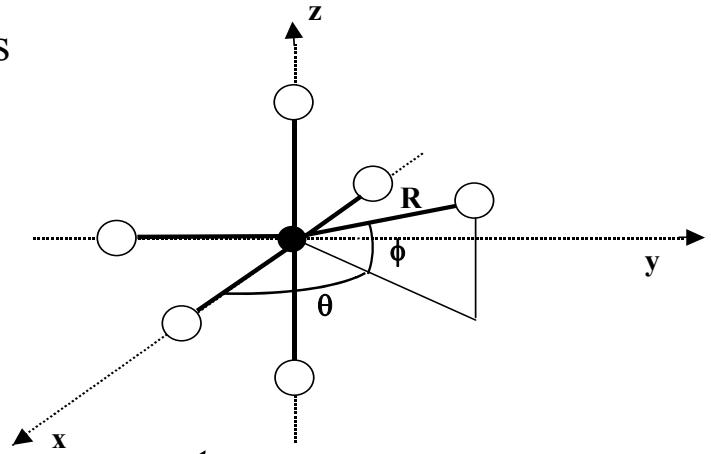
Exactly: all MS contributions are included

We have developed a new fitting method, called **MXAN**, that use the exact calculation of the scattering path operator

- i) We work in the energy space
- ii) We can start from the edge
- iii) We can use polarization dependent spectra

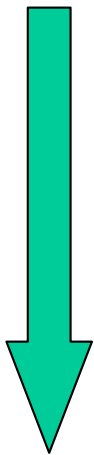
# To perform structural fits

- Initial geometrical configurations
- Exp. data



We generate hundred of theor. spectra by moving atomic coordinates

The potential is calculated at each step – Norman criterion



← Minimization of error function

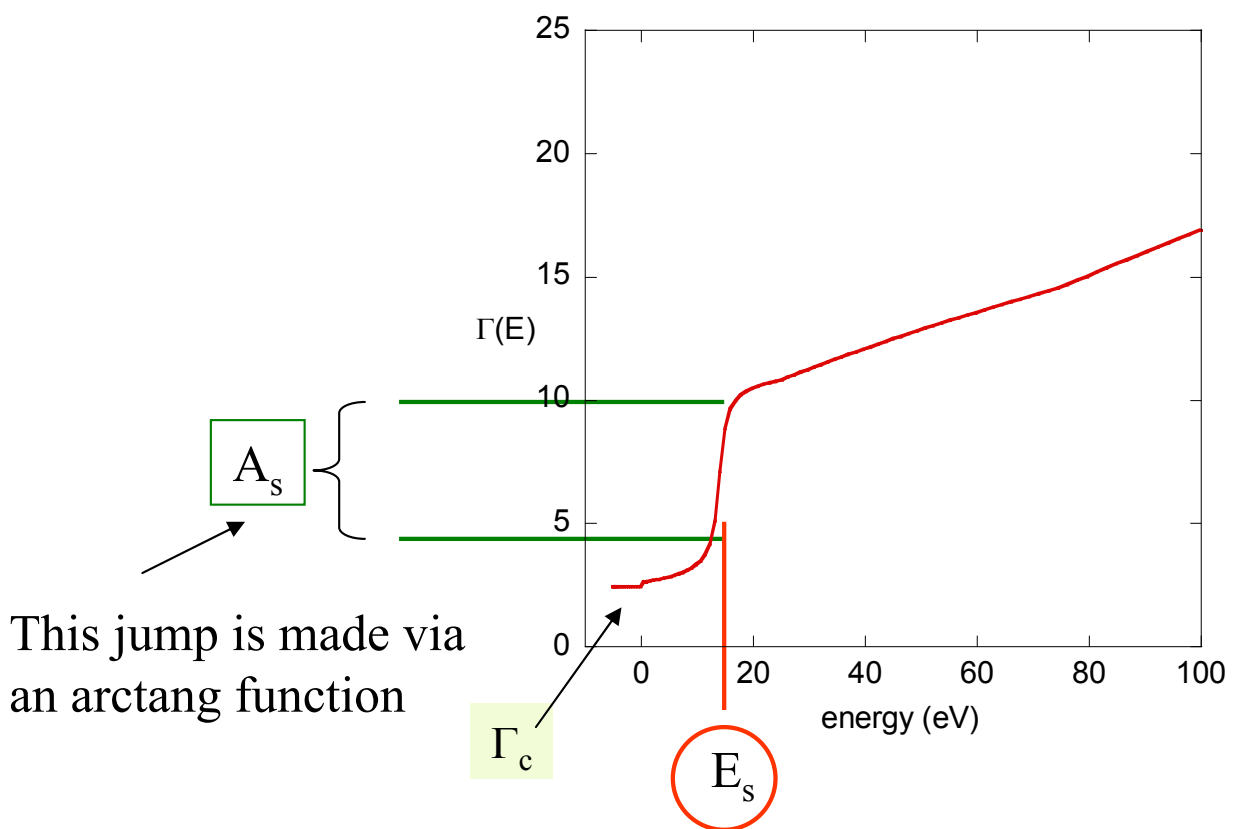
$$R_{sq}^2 = \frac{\sum_{i=1}^N \{ [y_i^{th.}(\dots, r_n, \theta_n, \dots) - y_i^{exp.}]^2 / \epsilon_i^2 \} w_i}{\sum_{i=1}^N w_i}$$

**By comparison with exp. data we can fit relevant structural parameters**

M. Benfatto and S. Della Longa (2001) J. Synch. Rad. 8, 1087  
 S. Della Longa et al. (2001) PRL 87, 155501  
 M. Benfatto et al. (2003) J. Synch. Rad. 10, 51

To avoid the problems coming from the use of the full HL potential, the Exchange and Correlation part of the potential is calculated on the basis of a phenomenological approach  $\longrightarrow$  real HL potential + convolution via a Lorentzian function with  $\Gamma_{\text{tot}}(E)=\Gamma_c+\Gamma(E)$

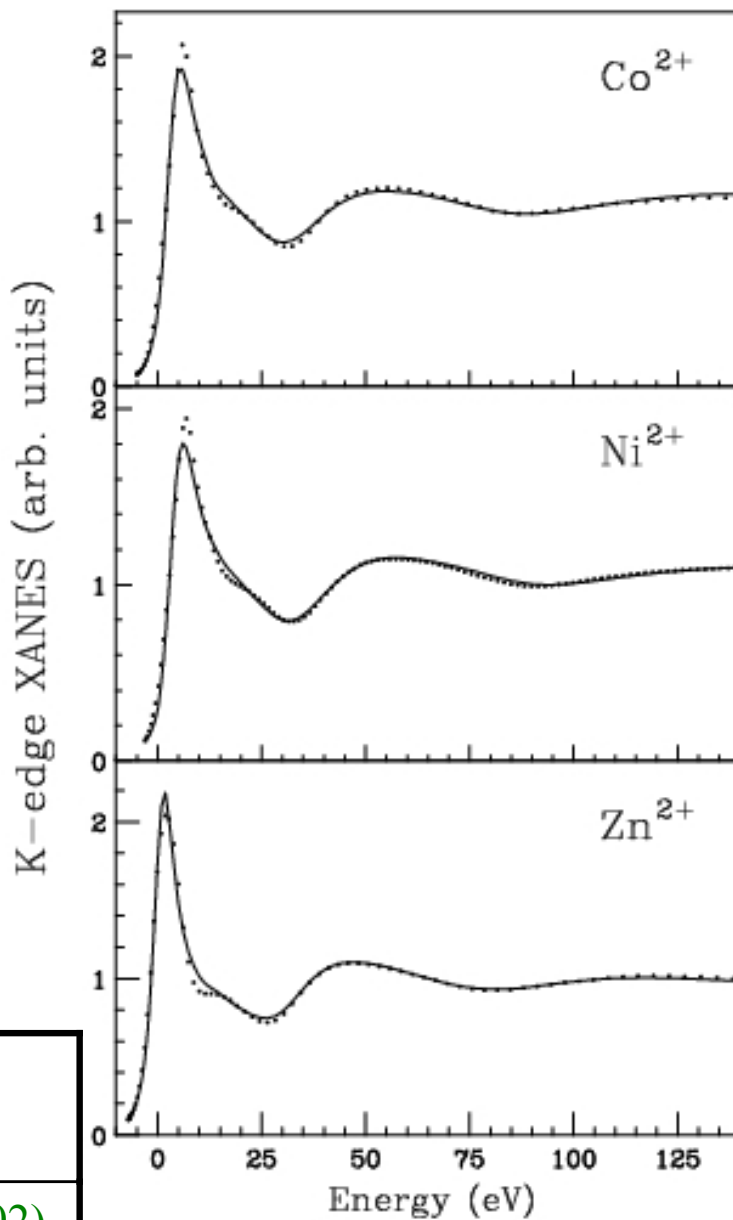
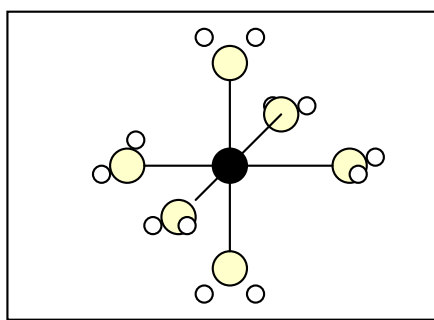
$\Gamma(E)$  Behaves like the universal form and starts from energy  $E_s$  with a jump  $A_s$ . Both  $E_s$ ,  $\Gamma_c$  and  $A_s$  are derived at each step of computation on the basis of Monte Carlo fit.



$\Gamma(E)$  contains all the intrinsic and extrinsic inelastic processes

# Transition metals in water solution

fits include Hydrogen atoms

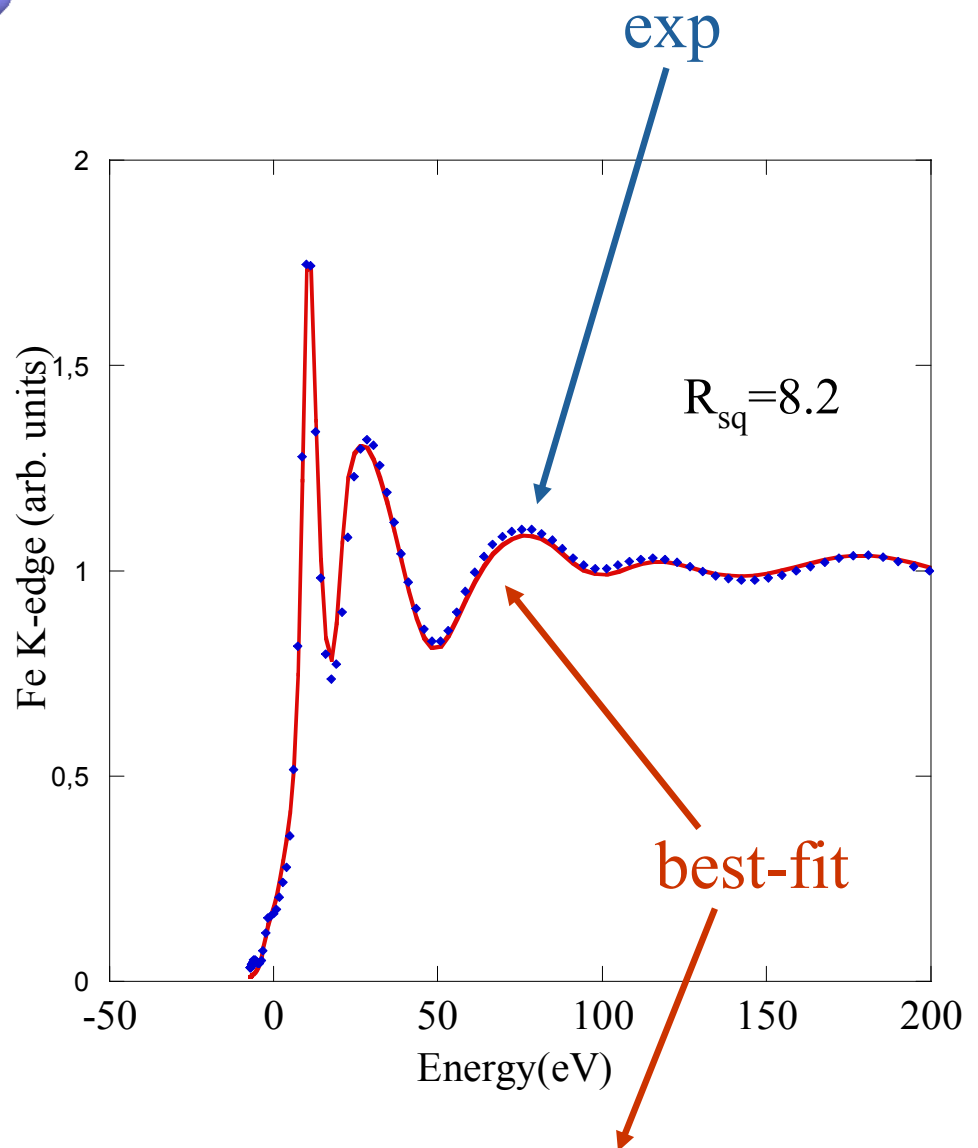
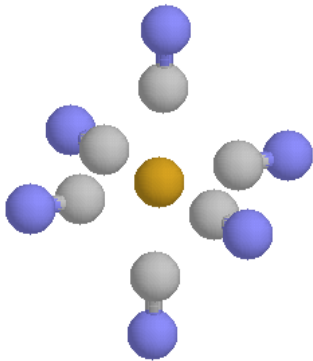


MXAN

GNXAS

	MXAN $R(\text{\AA})$	GNXAS $R(\text{\AA})$
$\text{Co}^{2+}$	2.06(0.03)	2.092(0.002)
$\text{Ni}^{2+}$	2.04(0.03)	2.072(0.002)
$\text{Zn}^{2+}$	2.06(0.02)	2.078(0.002)

# Fe (CN)<sub>6</sub> in water

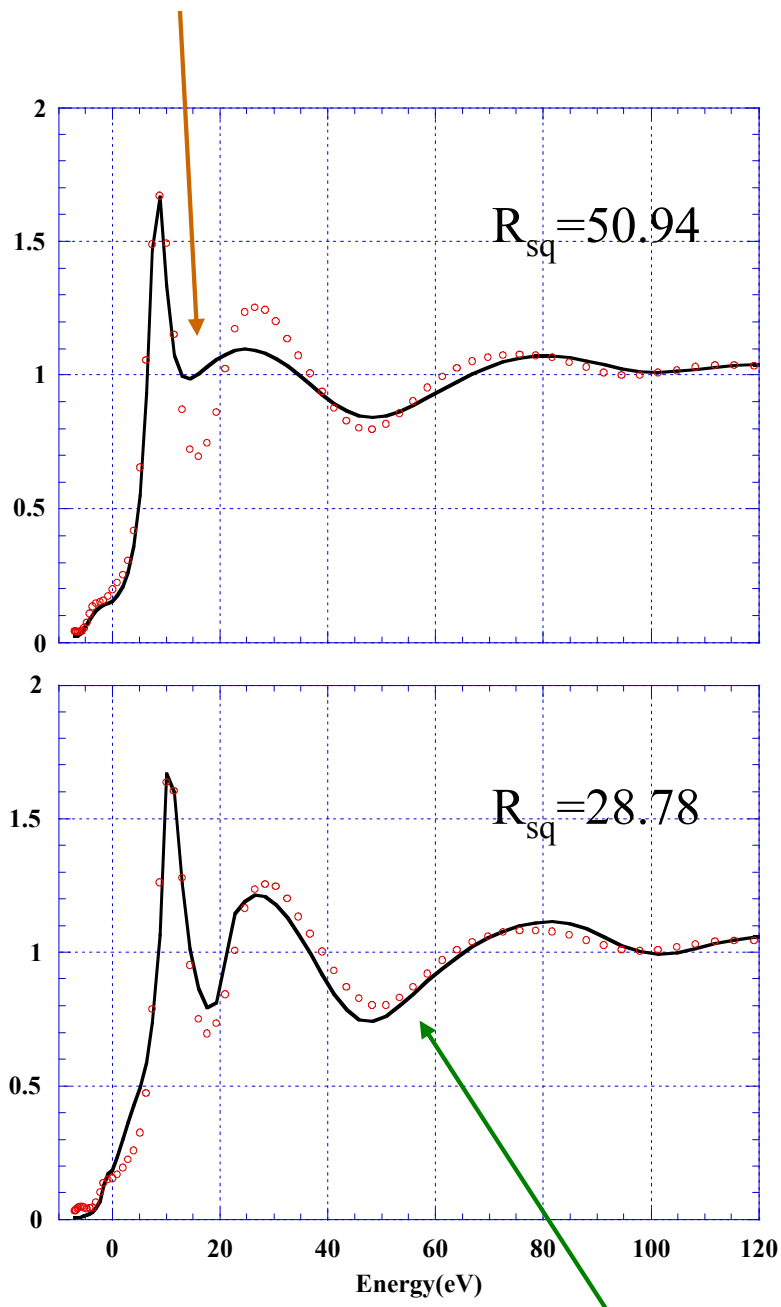


The best-fit condition corresponds to an octahedral symmetry with Fe-C distance of 1.92(0.01) Å and C-N distance of 1.21(0.01) Å

Previous GNXAS analysis (Westre et al. JACS 117 (1995)) reports Fe-C and Fe-N distances of 1.92 Å and 1.18 Å respectively

Fit with NO molecules

Fe-N = 1.91Å and N-O = 1.16Å



Fit with CO molecules

Fe-C = 1.94Å and C-O = 1.11Å

Chemical sensitivity



Test cases indicate a good structural reconstruction at the atomic resolution

high sensitivity to the structural changes

the fit results are weakly affected by the errors in the potential determination



just an increase of the error bar in the structural numbers

Three key points:

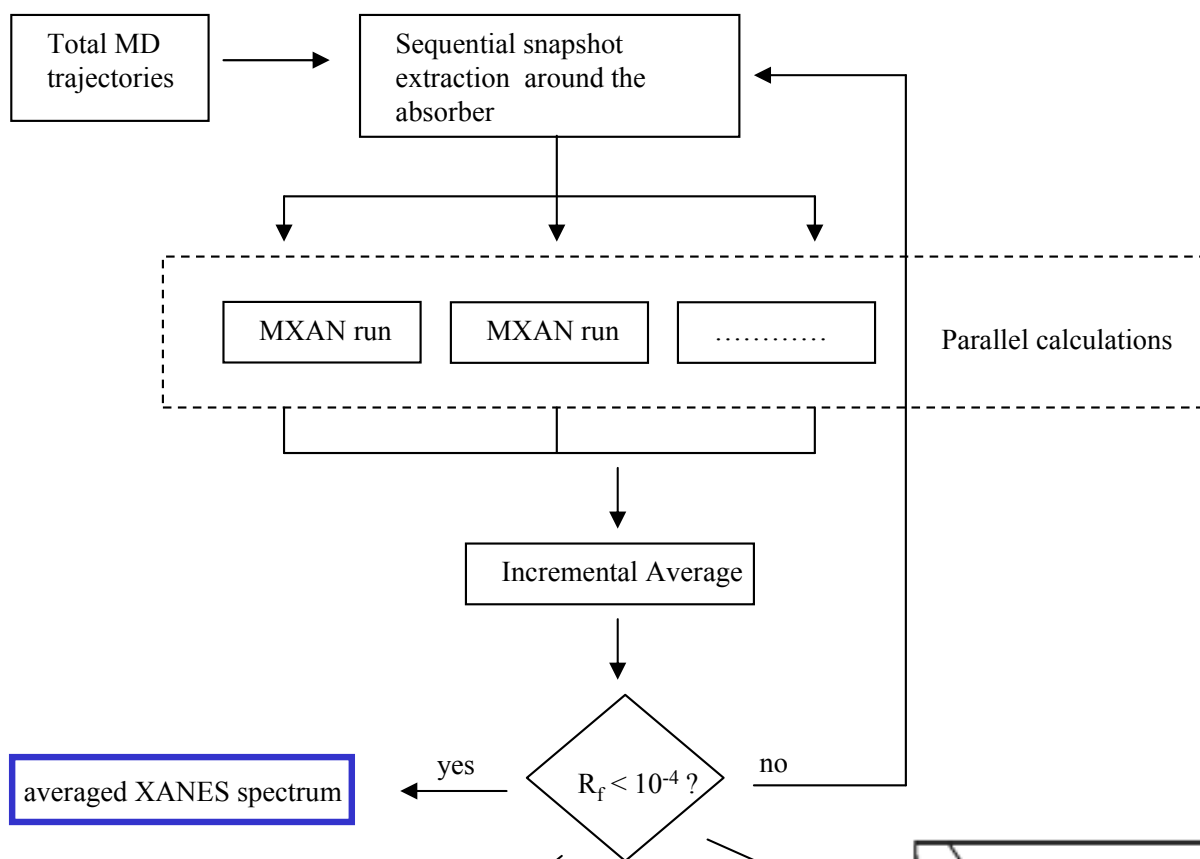
The use of a phenomenological damping

The potential is calculated at each step of the atomic movement keeping the same Norman criterion

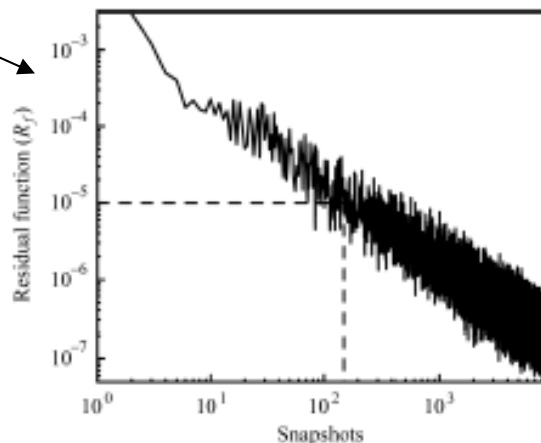
The used energy range is wide enough to minimize errors in the potential determination

# thermal and structural disorder

We use MD to generate thousands of geometrical configurations – each snapshot with a time step of 50 fs is used to generate one XANES spectrum – average using  $\sim 10^4$  geometrical configuration

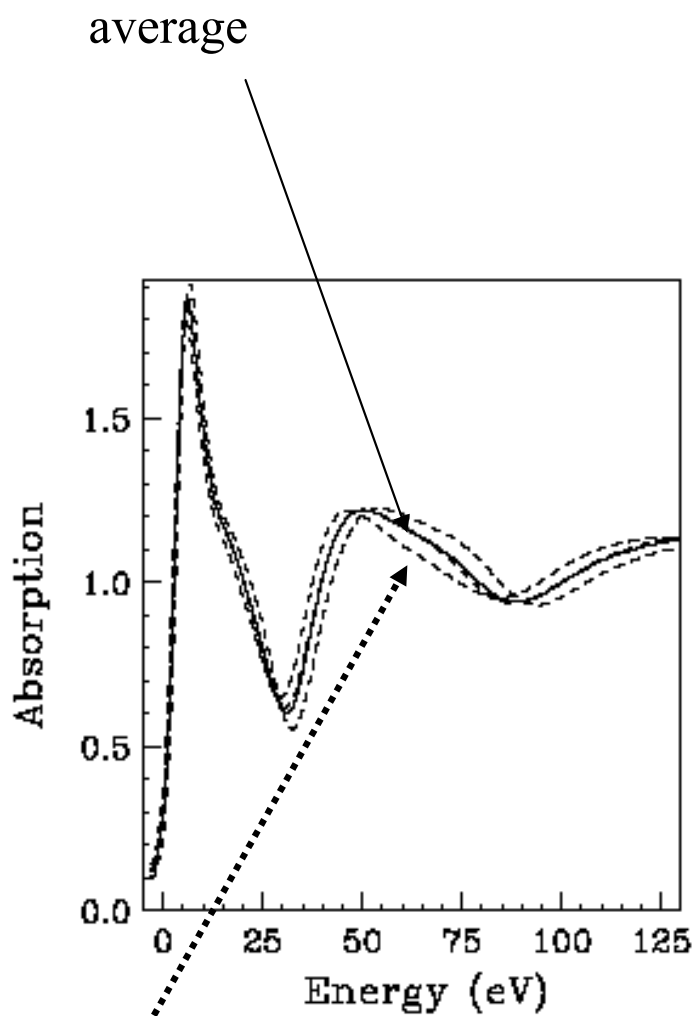
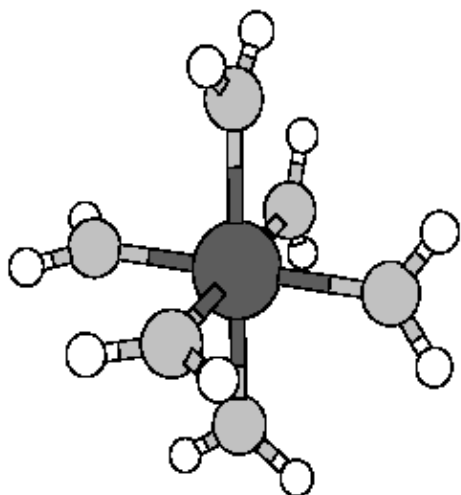


$$R_f(N) = \left[ \sum_i [\sigma^N(E_i) - \sigma^{N-1}(E_i)]^2 \right]^{1/2}$$



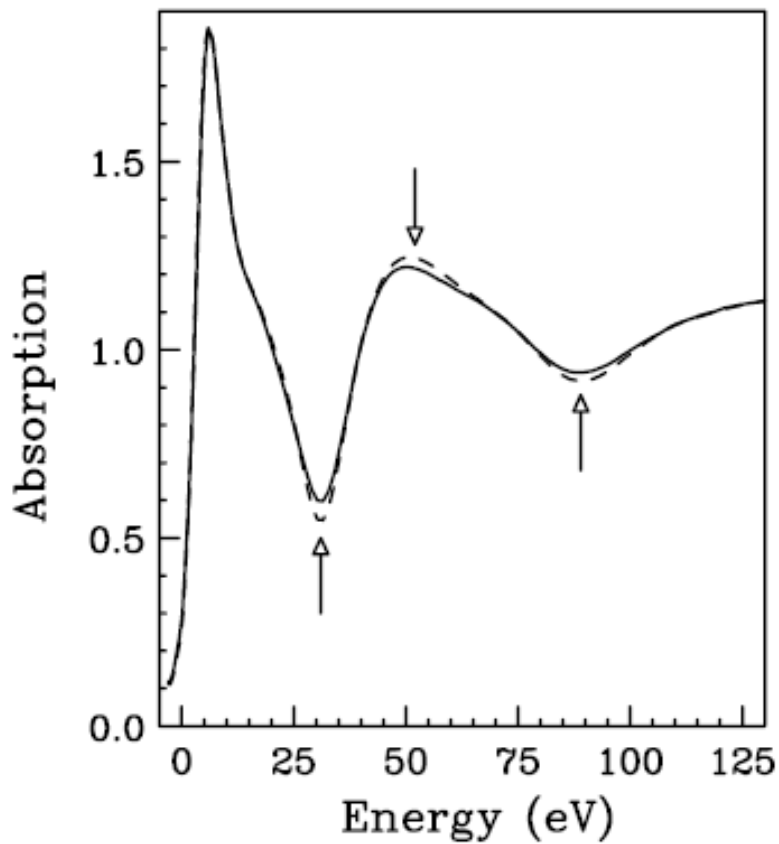
Merkling et al. JACS 124 (2002)  
Campbell et al. Jour.Synch.Rad. 6 (1999)

# Ni<sup>2+</sup> in water – Ni Kedge



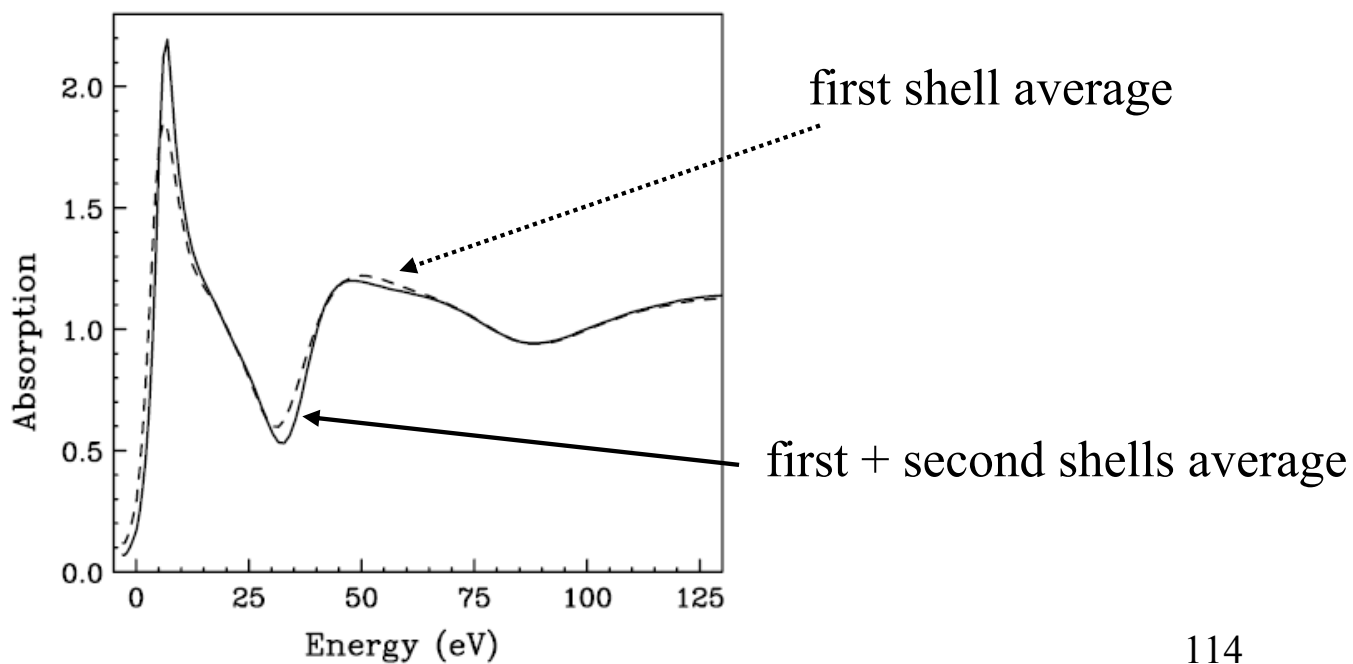
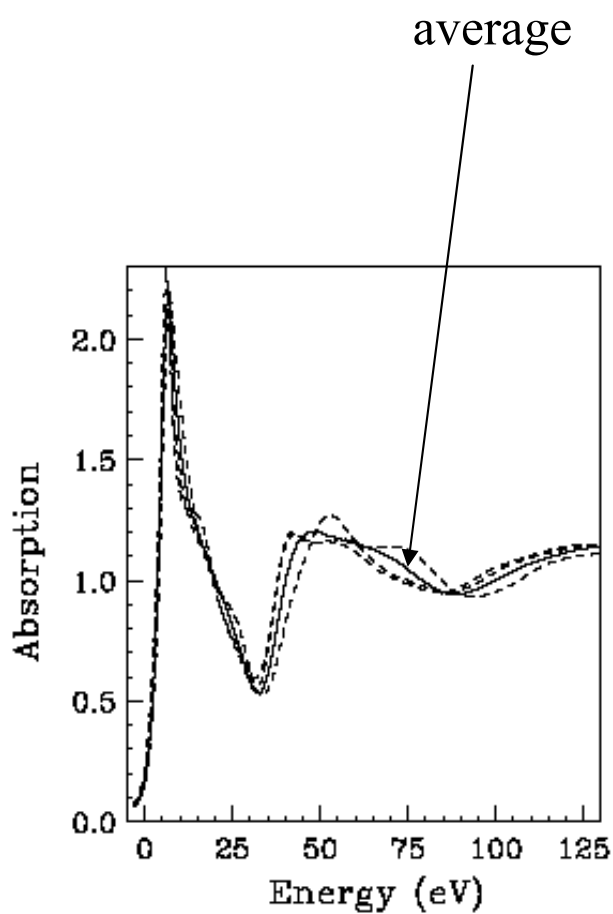
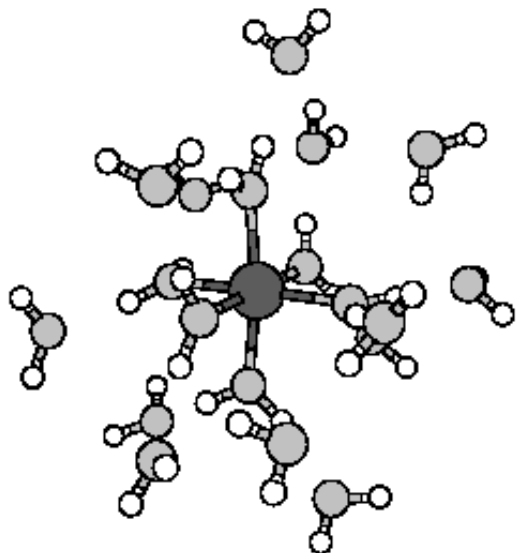
Calculations for some particular snapshots

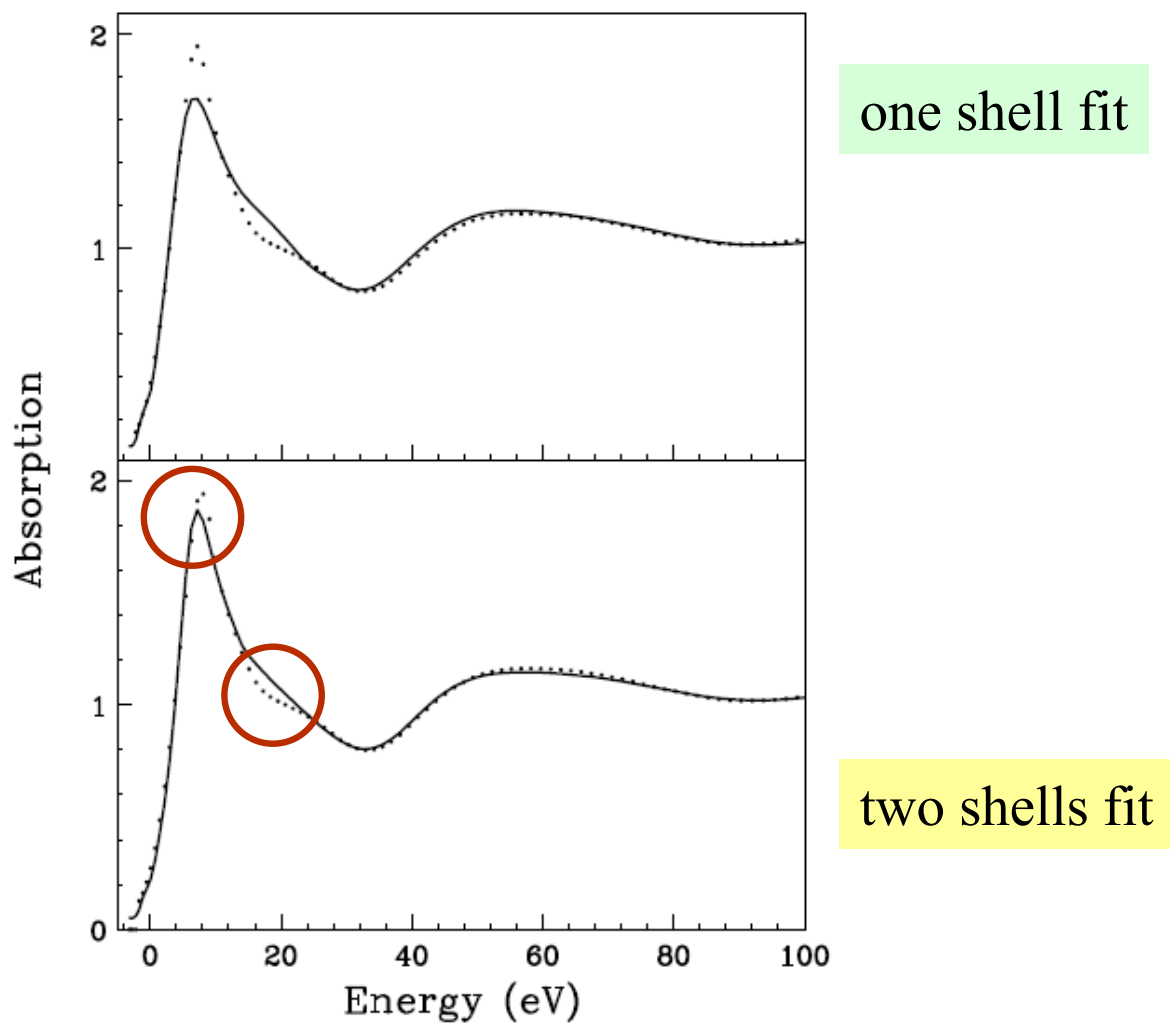
Comparison between the averaged theoretical spectrum and a single theoretical spectrum at the symmetrical first shell configuration



Arrows indicate the damping - very weak effect - it can be included in the phenomenological treatment of the inelastic losses in MXAN

including the second shell



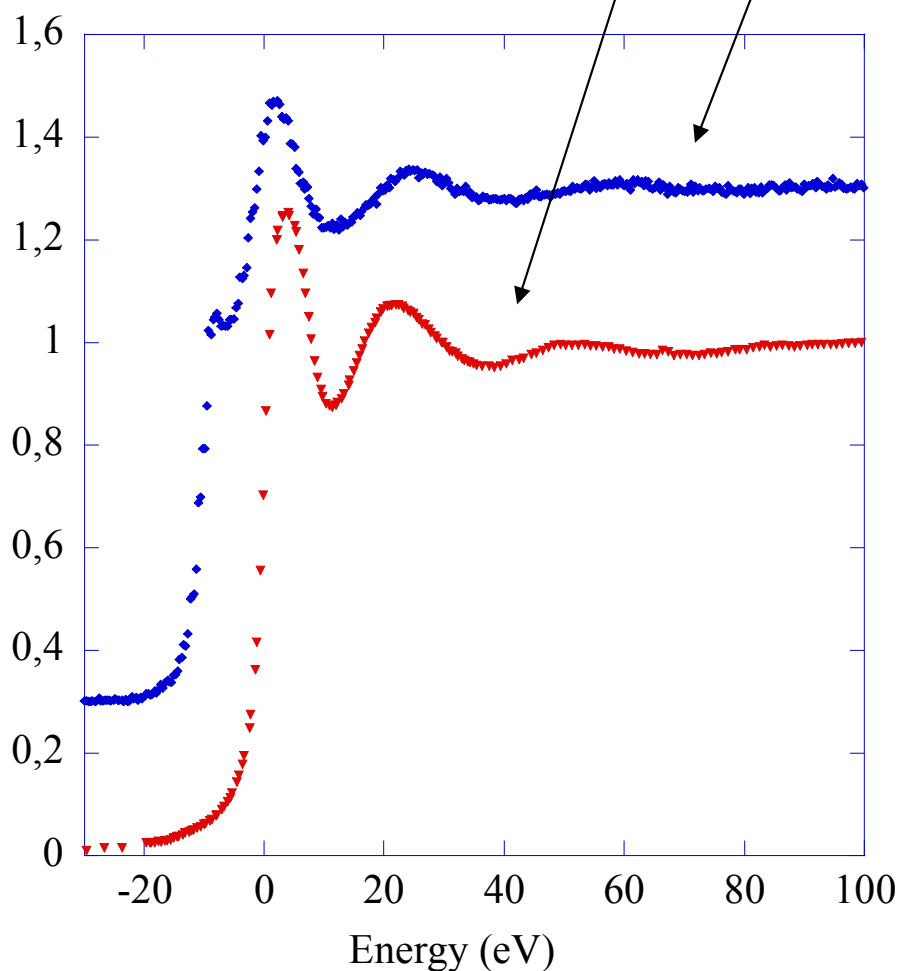


sizeable effects in the energy range 0 - 30 eV

P. D'Angelo et al. JACS 128 (2006)

## The water solvation of I<sup>-</sup>

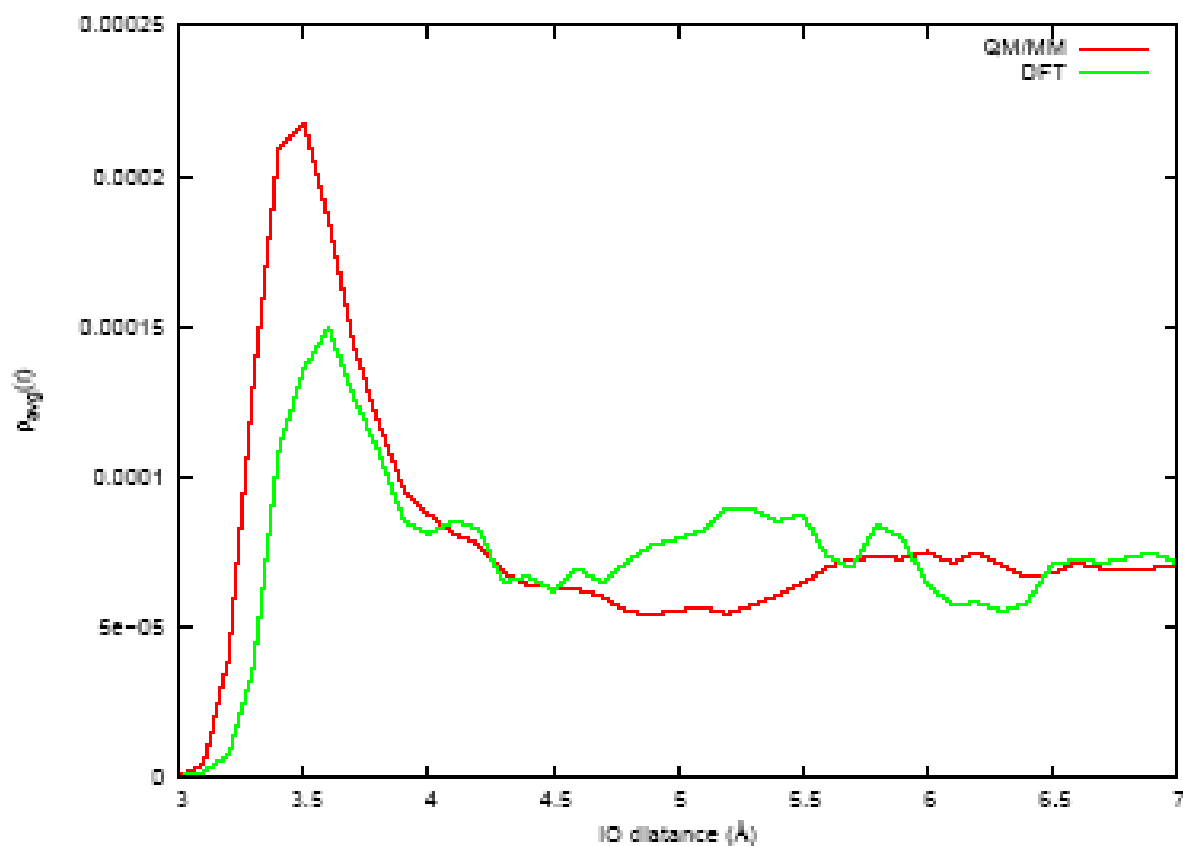
L1 - L3 XAS data



as in the previous case we analyze those data by MD snapshots generate by QM/MM and DFT methods

in collaboration with Chergui's group

## I-O pair distribution function from MD

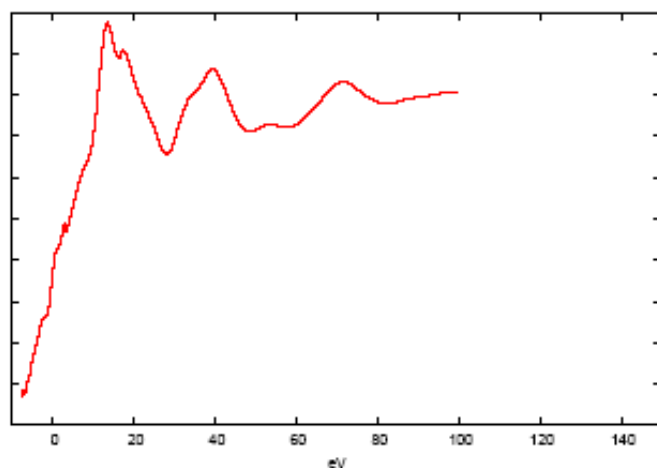
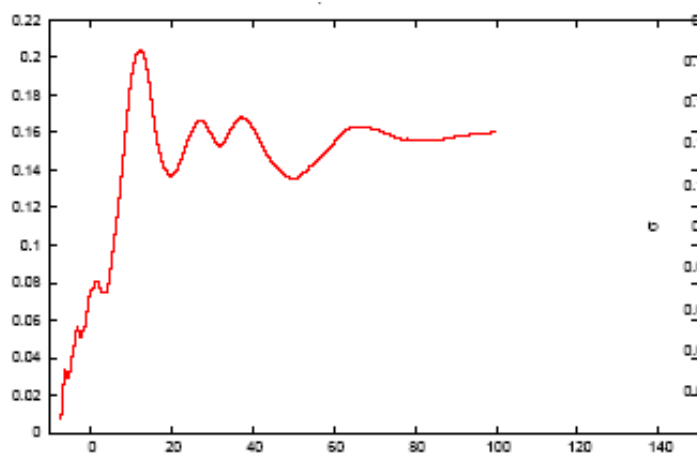
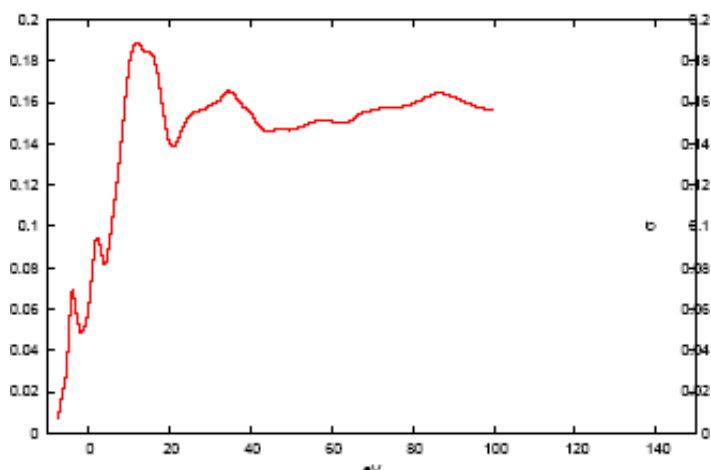


the calculated coordination number varies from 7 to 9.7 - difficulty to define a solvation shell



We have used more than 1000 frames - three of them at L3 edge

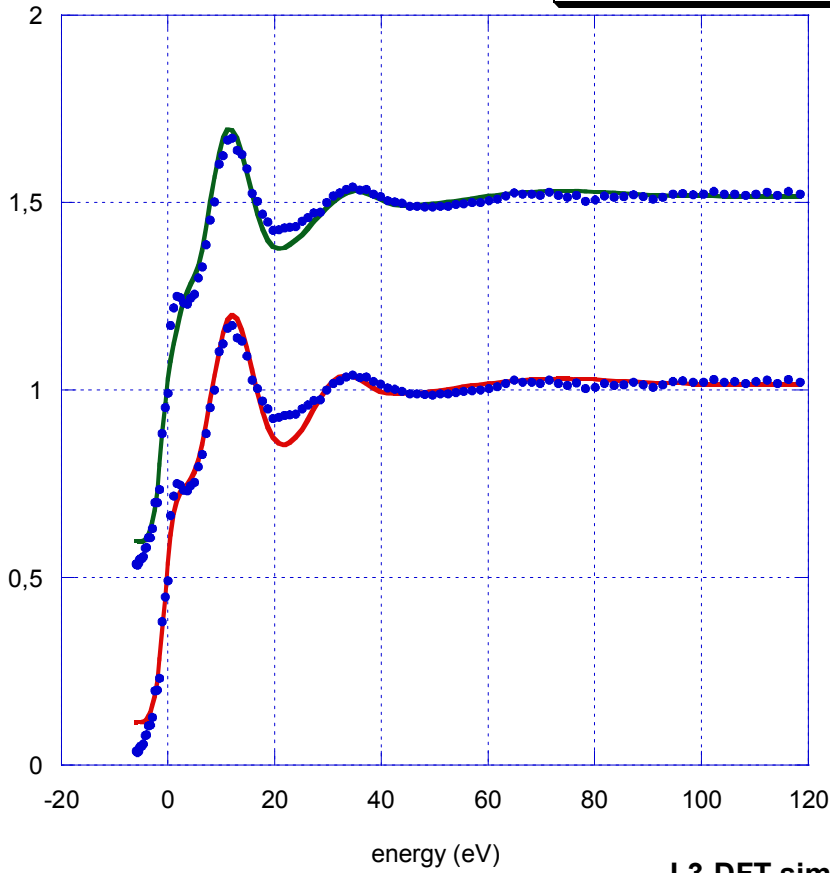
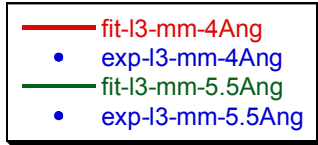
The calculation includes atoms (H and O) up to 7 Å



Very disordered system!

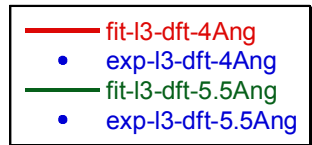
# L3 Fits

### L3-QM/MM simulation



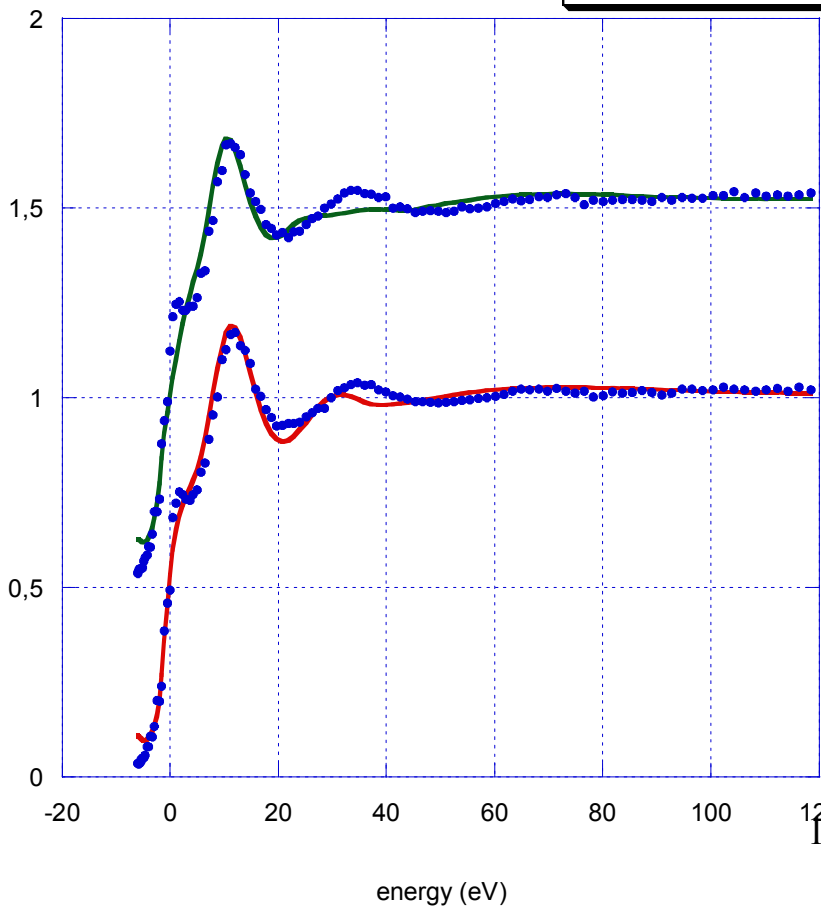
QM/MM

### L3-DFT simulation



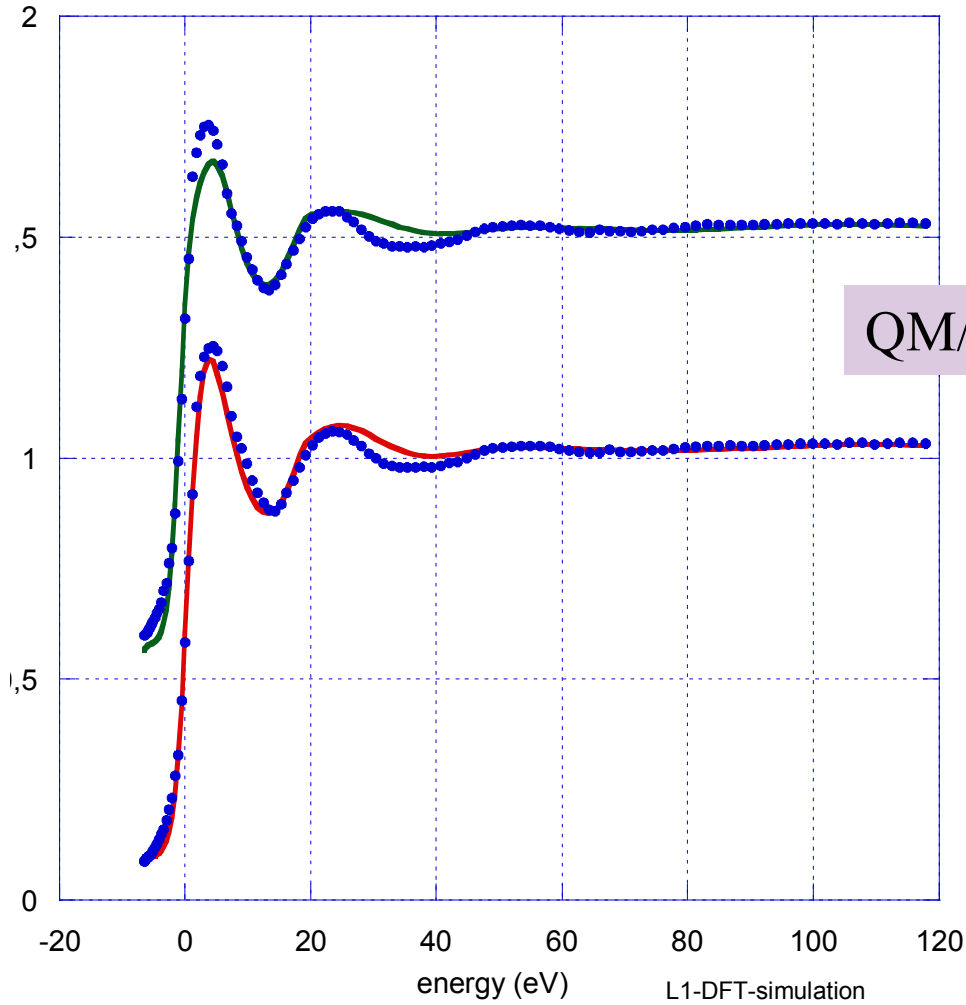
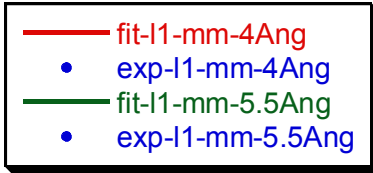
DFT

fit-I3-dft-4Ang

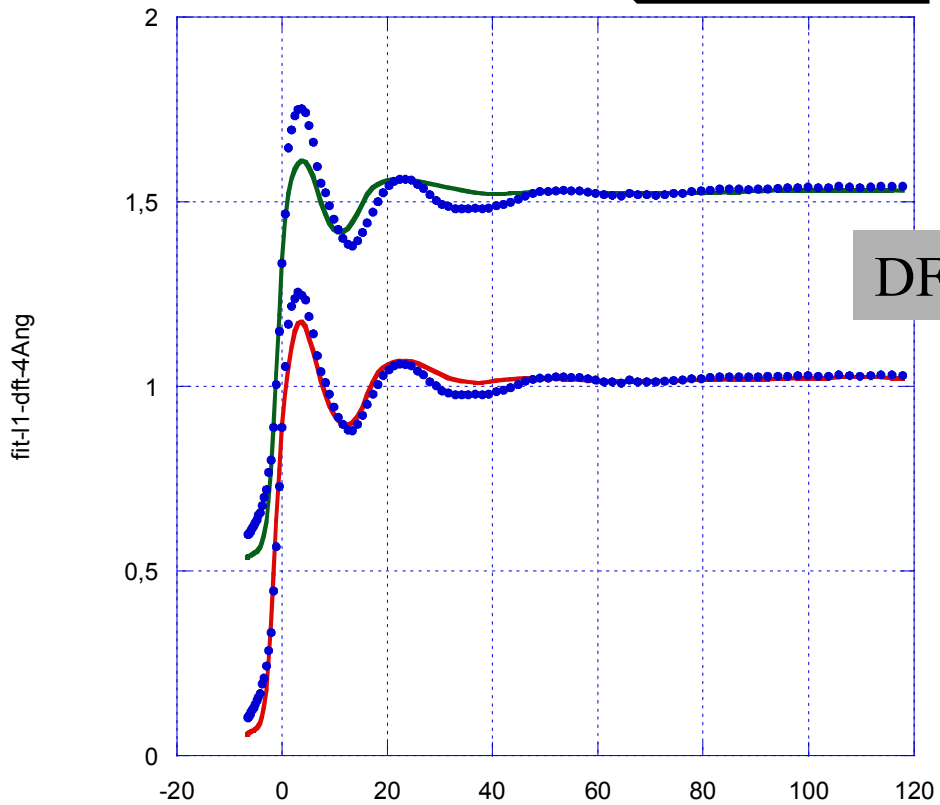
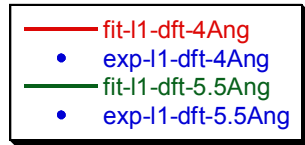


# L1 Fits

L1-QM/MM simulation

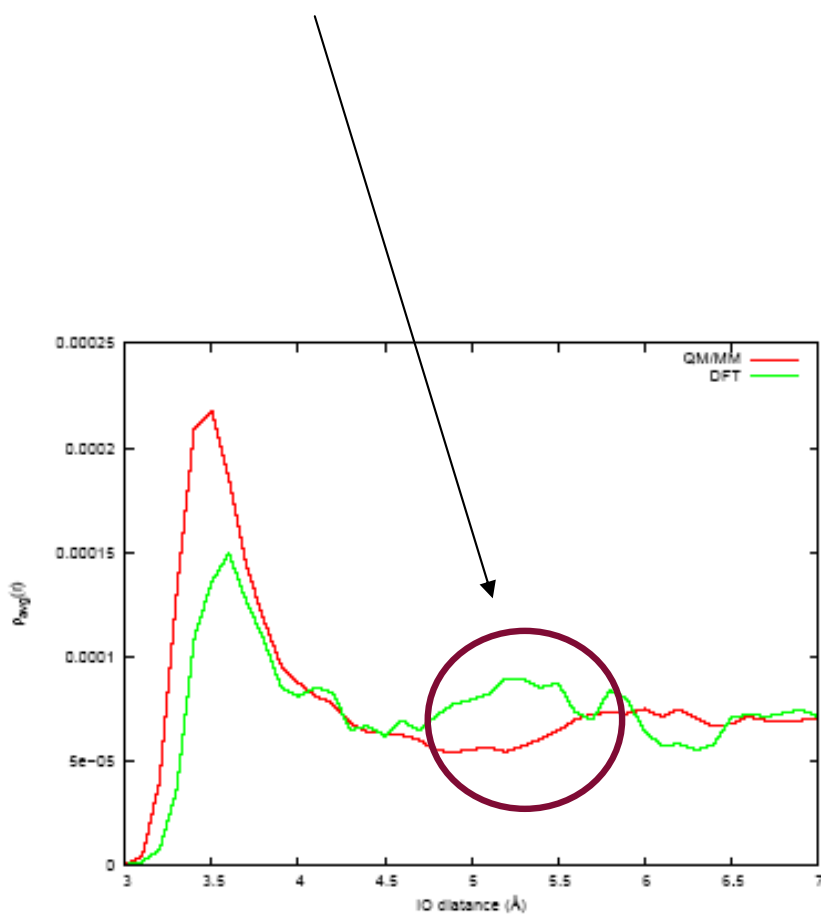


L1-DFT-simulation



The QM/MM calculations reproduces better than DFT the experimental data for both L1 and L3 edges - Increasing the cluster size DFT becomes worse than QM/MM

It seems that DFT introduces a partial order that is not verify in the reality



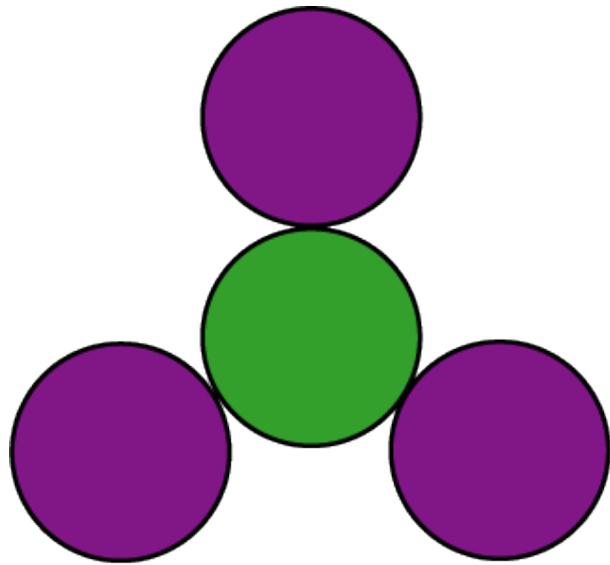
## some conclusions

It is possible to fit the XANES energy range starting from the edge to obtain quantitative structural information

We can treat, although numerically, very disorder systems – possible application on nano-materials

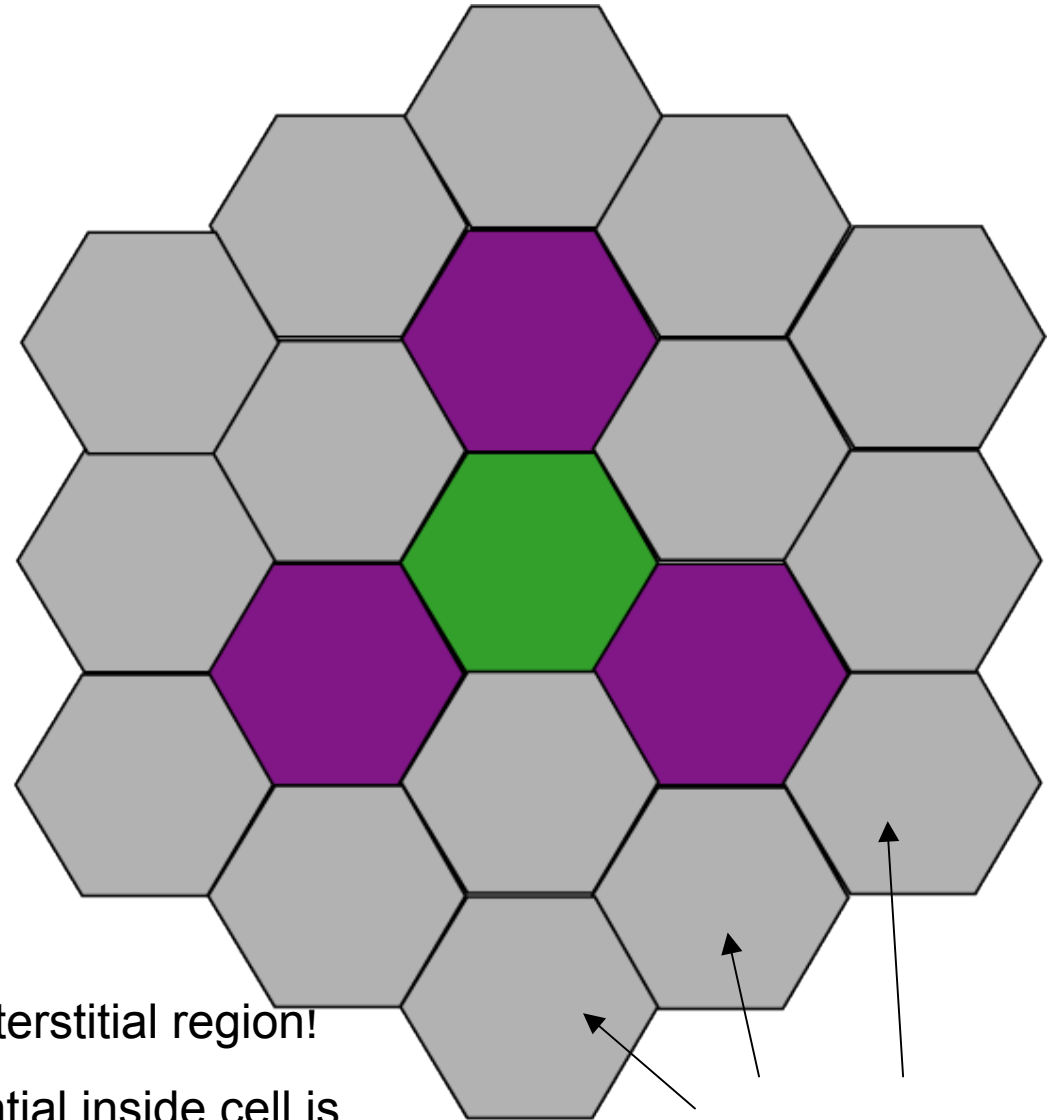
Future: new MXAN using the non-MT theory  
- improvements for the exchange-correlation potential

### MT approximation (MXAN)



Use spherical shaped and averaged potential  
outside the cells potential is flat

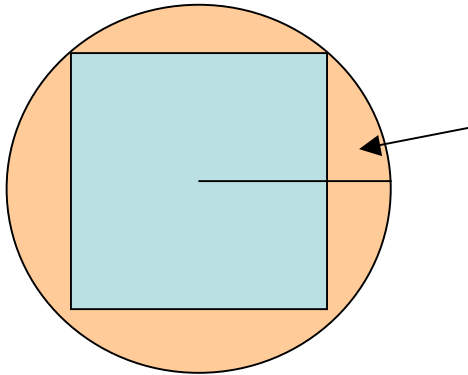
### Non-MT (FPMS)



No interstitial region!  
Potential inside cell is anisotropic

Empty Cells (EC)

# FPMS

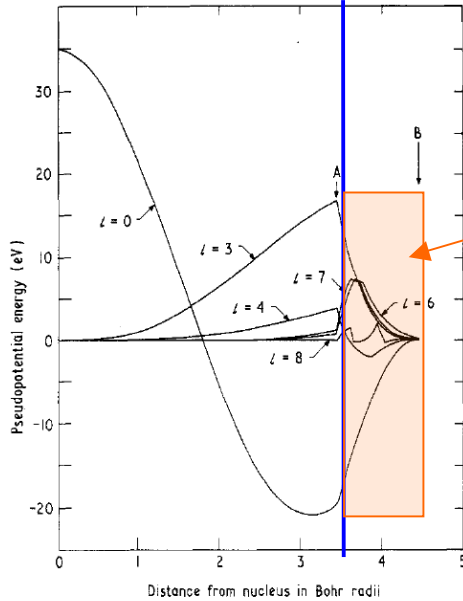


So called moon region.  
In this region  
expansion by  
spherical harmonics  
never converges.

There are two types of  $I_{\max}$  in MS

1. potential power
2. anisotropy

$$V(\mathbf{r}) = \sum_L \underline{V_L(r)} Y_L(\hat{\mathbf{r}}) \quad (L=(l,m))$$



For discontinued function  
spherical harmonics  
expansion  
**never converges** in  
practice.  
(Gibbs phenomenon)

Radial component oscillate!

Figure 8. Spherical harmonic components of the pseudopotential of Cohen and Bergstresser. A, muffin tin radius; B, radius of circumscribed sphere.

=> We expand the wave function.

$$\Phi(\mathbf{r}) = \sum_L R_L(r) Y_L(\hat{\mathbf{r}}) \quad (L=(l,m))$$

Since the wave function and its first derivative are continuous, the series converges uniformly even with truncated potential.

Kellog, *potential theory* (1967)

$$L^2\Phi(\mathbf{r}) = \sum_L l(l+1)R_L(r)Y_L(\hat{\mathbf{r}})$$

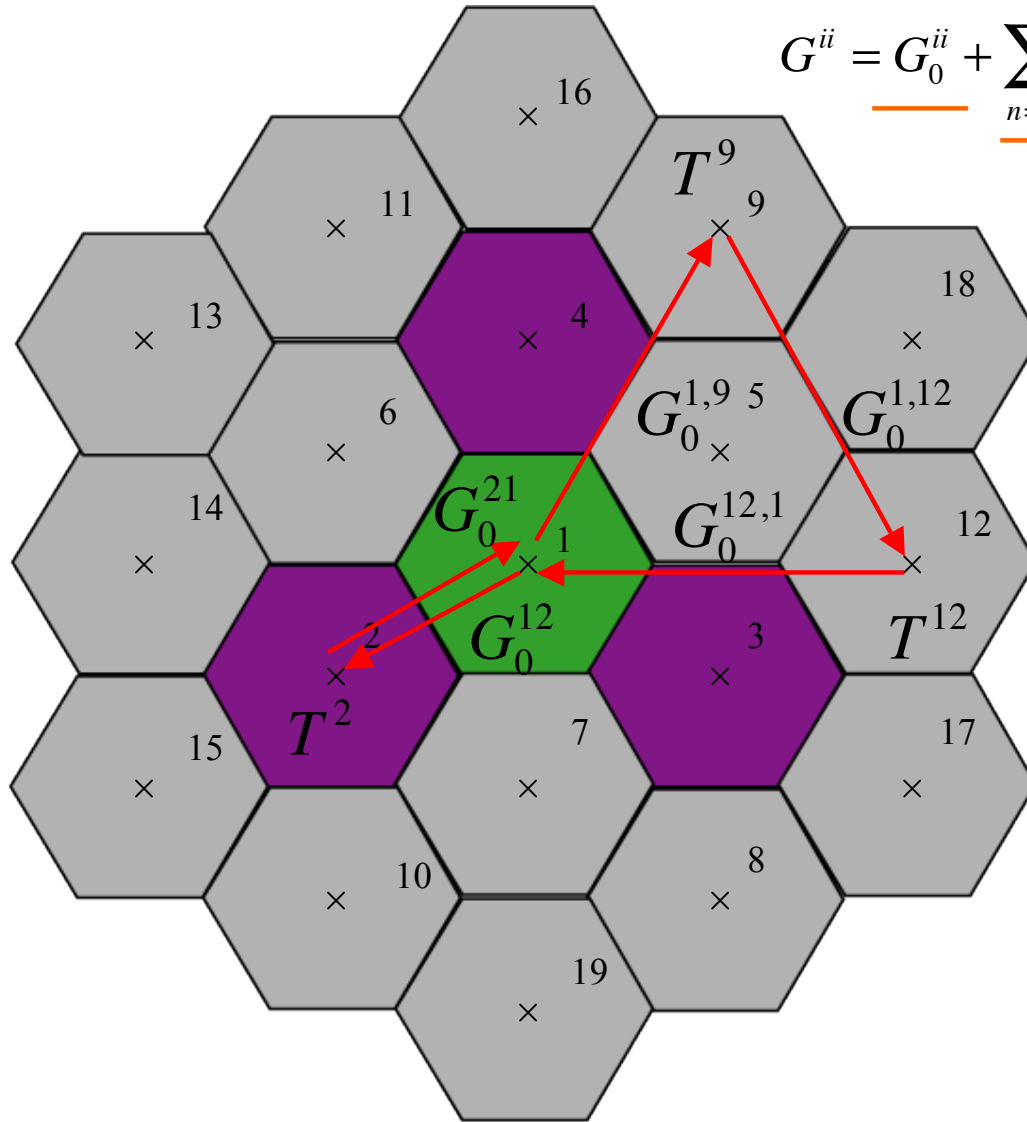
We treat it as inhomogeneous term in Schrödinger equation.

To solve the equation efficiently, we developed modified Numerov method.

K. Hatada et al, PRB (2007)



## Multiple scattering : Green's function picture



$$G^{ii} = \underbrace{G_0^{ii}} + \underbrace{\sum_{n \neq i} G_0^{in} T^n G_0^{ni}} + \underbrace{\sum_{n \neq m \neq i} G_0^{in} T^n G_0^{nm} T^m G_0^{mi}} + \dots$$

All sites have centers.

$G^{ii}$

: Green's function on site  $i$

$G_0^{ii}$  : isolated

$$\sum_n G_0^{in} T^n G_0^{ni}$$

: single scattering

$$\sum_{nm} G_0^{in} T^n G_0^{nm} T^m G_0^{mi}$$

: double scattering

**First-principles calculation of x-ray absorption-edge structure in molecular clusters**

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*Department of Applied Physics and Stanford Synchrotron Radiation Laboratory, Stanford University, Stanford, California 94305*

F. W. Kutzler

*Department of Chemistry, Stanford University, Stanford, California 94305*

(Received 30 July 1979)

We report initial results of a systematic study of the calculation of near-edge features in the x-ray absorption spectra of an atom in a molecular cluster as a function of different chemical environments and of varying prescriptions for the effective molecular potential of the excited electron. For a test-case comparison of the  $K$  edges of  $\text{GeCl}_4$  and  $\text{GeH}_4$ , we find that the chemically induced changes in the molecular potential lead to large changes in the occurrence and strengths of bound-state and shape-resonance spectral features, in semiquantitative agreement with observation.

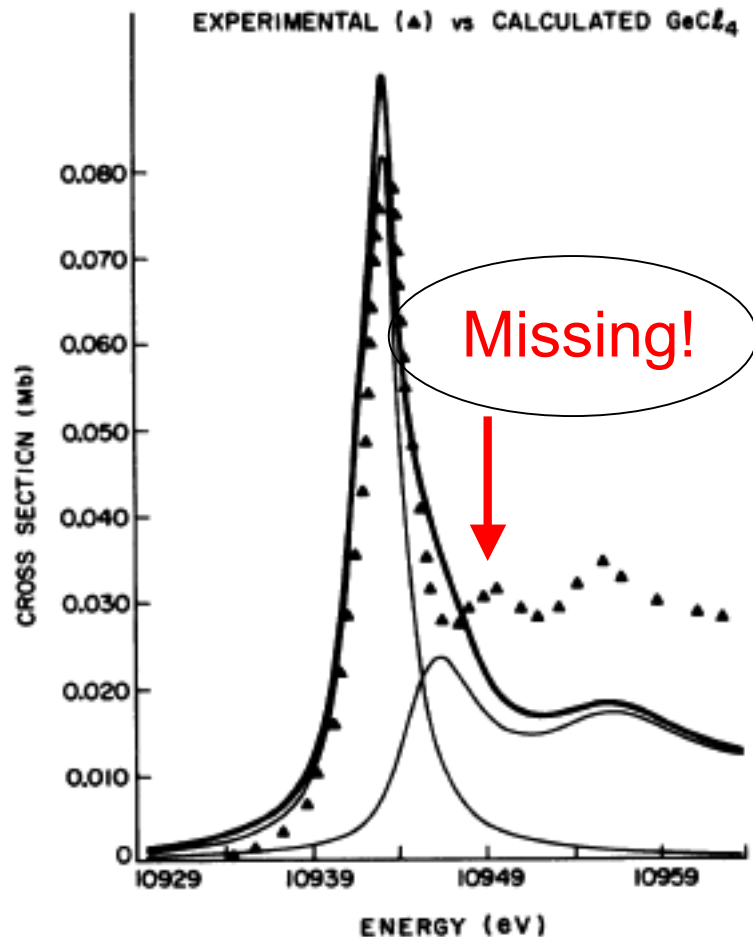
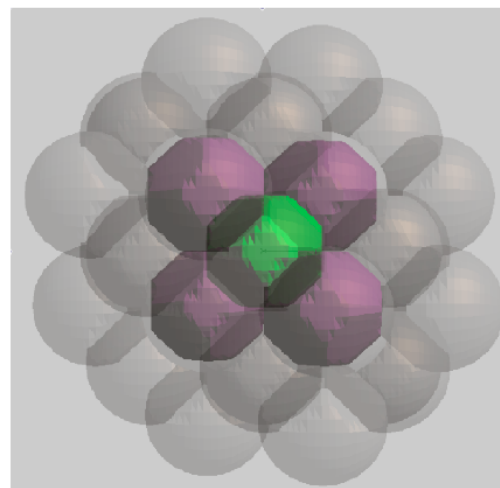
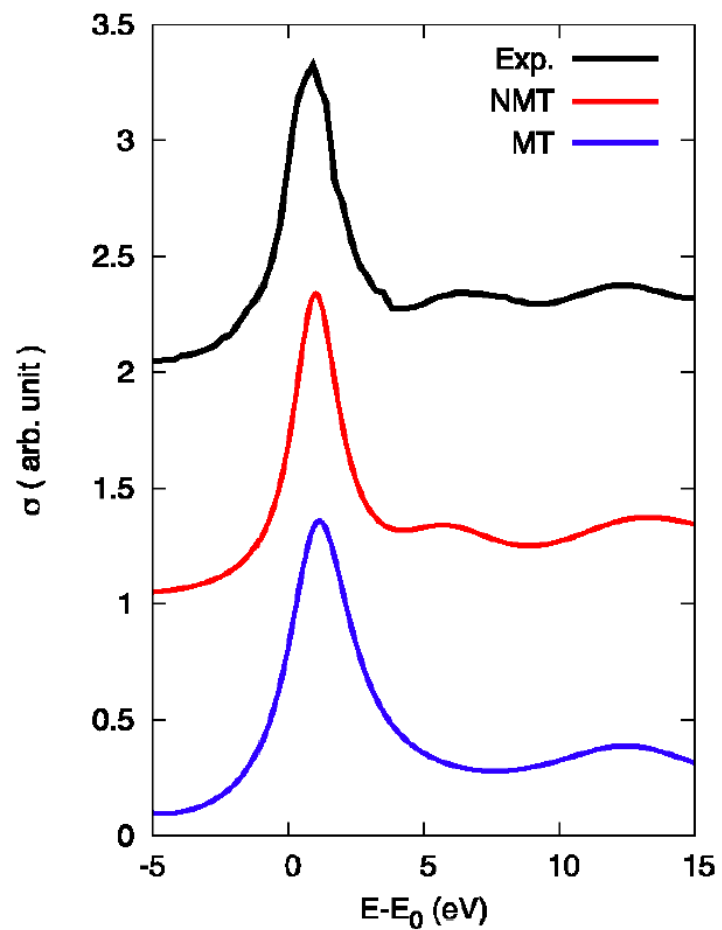


FIG. 1. Computed  $K$ -absorption spectra for Ge in  $\text{GeH}_4$  (upper panel) and  $\text{GeCl}_4$  (lower panel) after convolution with a Lorentzian to represent lifetime and experimental resolution effects. The light curves represent the bound-state and continuum contributions separately, while the heavy curves give the sum of the two. The energy placement of the experimental points has been adjusted for a reasonable fit (see text). Note that inelastic effects will generally be expected to increase the total cross section over the results given by the present calculation.

As far as energy dependence goes, the overall vertical scale of the experimental points shown in Fig. 1 for  $\text{GeCl}_4$  has been adjusted by eye to give an overall "best fit." It may be noted that although the number and magnitude of the calculated spectral features are in qualitative agreement with experiment, the separation of the first and second continuum resonances is of order 10.0 eV compared to 6 eV seen experimentally. We ascribe this discrepancy to the inaccurate treatment of the interstitial potential by the use of the muffin-tin approximation.

Check Natoli's prediction of 30 years ago.

# Ge K-edge of GeCl<sub>4</sub>



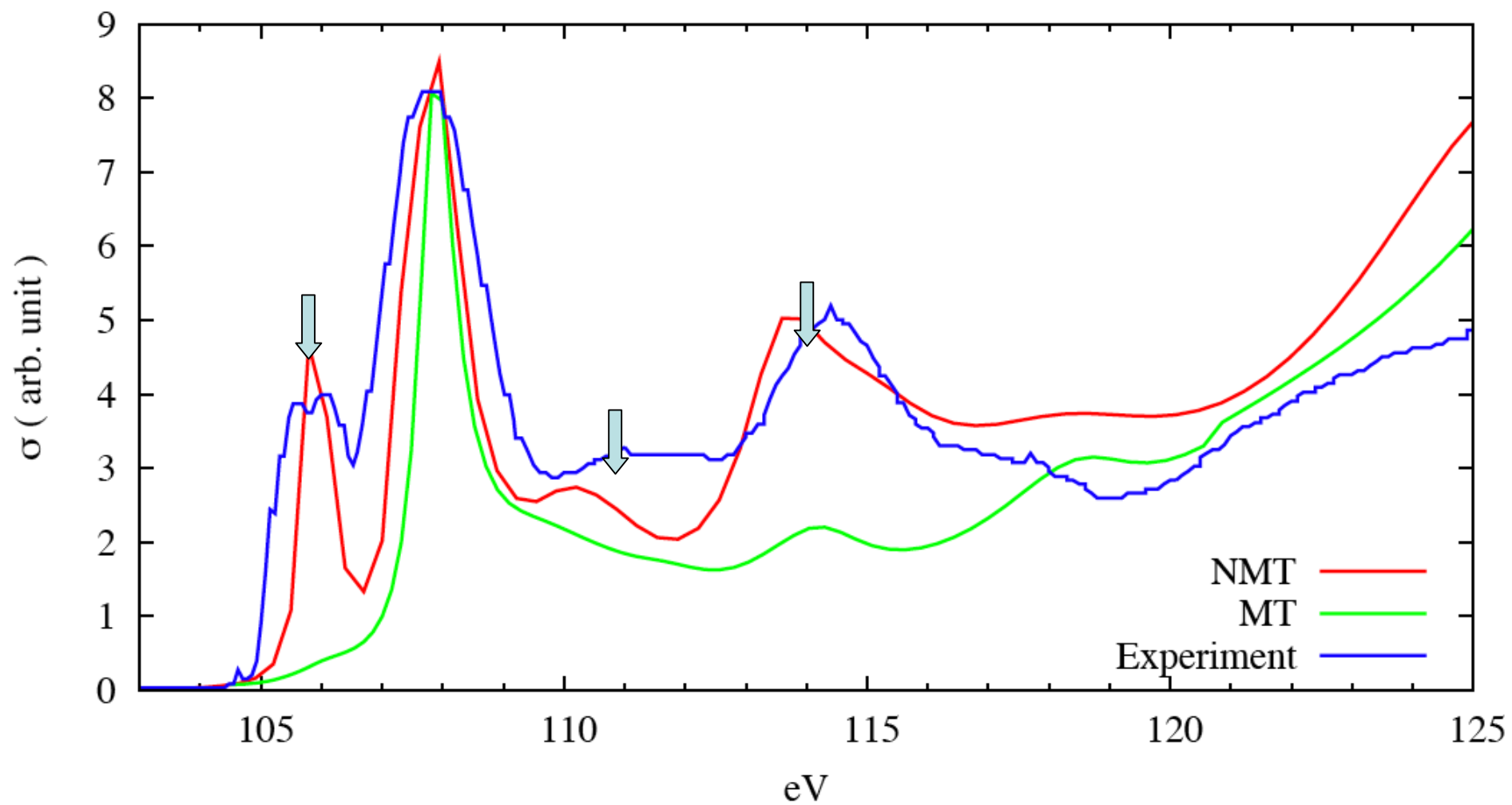
1 Ge + 4 Cl + 38 EC  
(Empty cells)

# Other spectroscopies

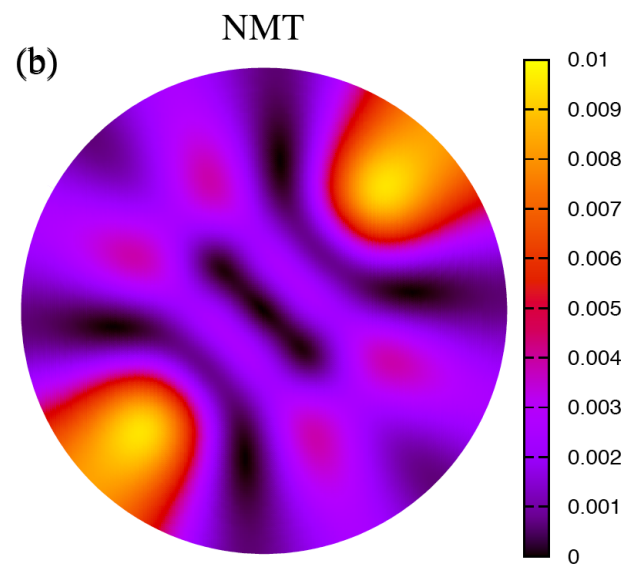
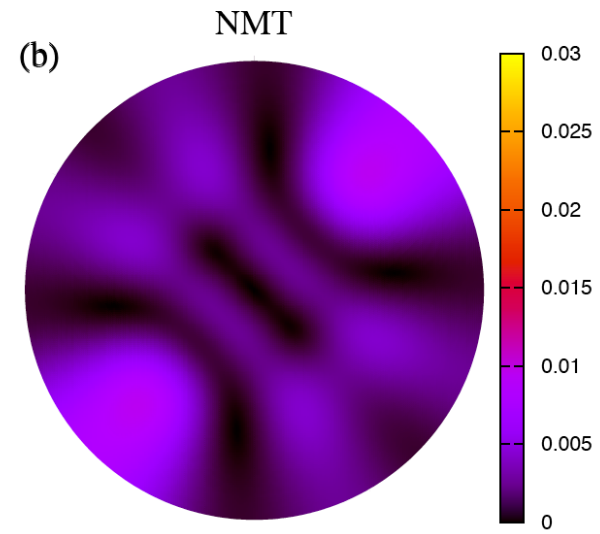
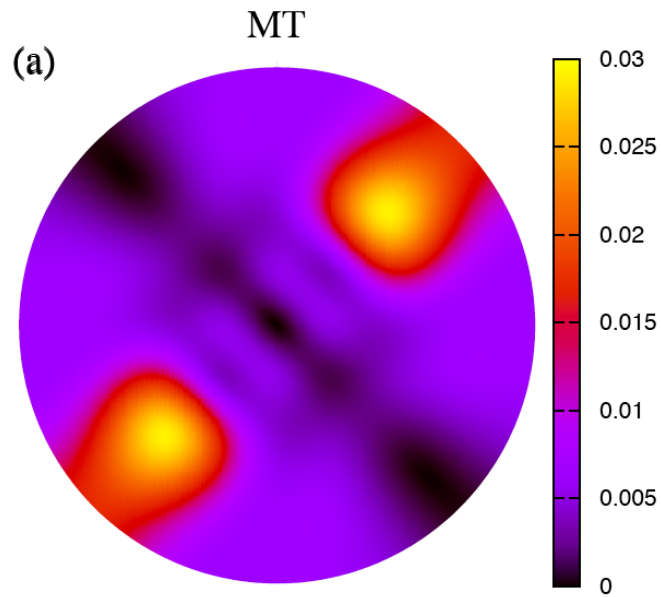
- Energy-Loss Near-Edge Structure (ELNES)
- Resonant X-ray Scattering (RXS)
- X-ray Photoelectron Diffraction (XPD)

# ELNES (NMT vs MT)

SiO<sub>2</sub> ( $\alpha$ -quartz) of Si-L<sub>2,3</sub>-edge with ( 5 Å cluster )



# PHD of Ge 1s of GeCl4



- You can find MXAN and FPMS programs at,

<http://www.Inf.infn.it/theory/CondensedMatter/index.html>

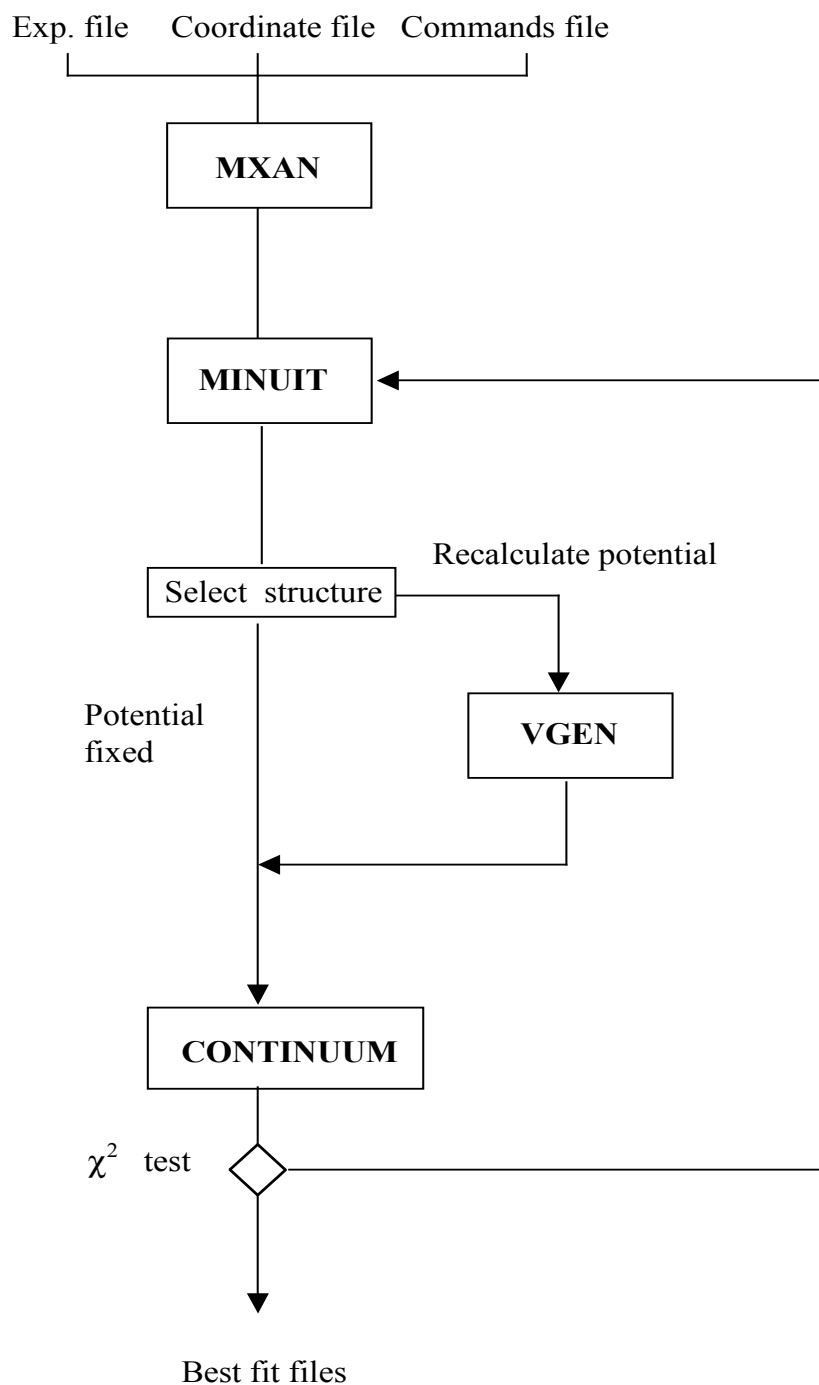
Programs are free, but registration is obligated.



# THE MXAN PROCEDURE

# How to run

## Flow diagram of MXAN



We need three inputs

**Experimental data:** two columns file with the normalized experimental data

```
energy(eV) abs
-10.00    0.034048
-9.000    0.037412
-8.000    0.043628
-7.000    0.052225
-6.000    0.068452
-5.000    0.098043
-4.000    0.14981
-3.000    0.23317
-2.000    0.35128
-1.000    0.53214
0.0000    0.74763
1.0000    1.0009
2.0000    1.2878
3.0000    1.5565
.....    .....
```

We can use any energy step

**Starting geometry file** – The prototype is generated by MXAN – Users must change it for own use

```
&JOB
fname ='NiH2'
nsca  = 19
lmax  = 3
absorber= 1
outersph=.false.
edge  ='k'
emin  = -3.000
emax  = 150.000
delta = 0.200
xscale = 1.400
gamma = .000
typot ='hrel'
charelx ='ex'
norman ='extrad'
ovlpfac = .120
coor  ='angs'
potcalc ='Y'
symcalc ='Y'
parcalc ='Y'
v0imp  = .000
rhoimp  = .000
&END
```

```
tit1
tit2
tit3
```

```
tit4 iz   x       y       z   rad  ilig ----- comm -----
```

```
28 .00000 .00000 .00000 1.26000  0 .00000 Ni_____000
 8 .00000 -1.90711 .00000 .90000  1 2.02119 O_____eq1
```

.....

-----  
---- CuN square piramid with O ax with H and more N --  
-----

CuNH --parameter- in.deviat. - init.err. --- phys.limits ---

1	r001	0.0000	0.0200	-0.1000	0.1000
2	f001	0.0000	3.0000	-20.0000	20.0000
3	t001	0.0000	0.0000	-10.0000	10.0000
4	r002	0.0000	0.0200	-0.1000	0.1000

.....

34	r012	0.0000	0.0000	-0.1500	0.1500
35	f012	0.0000	0.0000	-10.0000	10.0000
36	t012	0.0000	0.0000	-10.0000	10.0000

----- Control options -----

36 ! n of parameters (fixed + variable)  
NORM 1SHE ! Mode(NORM,EXTE,RECO),Search(1SHE,CXYZ,OPTN)  
U ! U (Unpol) XY (PolarXY) Z (Pol.z) P (Pol.)  
1 ! not active...

----- INPUT files -----

CuNH3.exp ! EXP. FILE to fit  
0.01000 ! EXP. ERROR  
DATA.CuN8OH ! Coords and options for potential  
N 1 ! Other th. calc. sites? How many?  
0.5 CuN2H-22.BEST.0.0 ! Fract,Input calc.othersite, Shift?

----- INTERMEDIATE files -----

CuNH.PARM ! params for MS  
CuNH.SYMM ! File with symmetry info and lmax  
CuNH.XALP ! X-alpha potential  
CuNH.COUL ! Coulomb potential  
CuNH.RHOD ! Charge densities

----- OUTPUT PREFIX -----

CuN8OH-1 ! Output files prefix (MAX 8 CHAR)

----- PARAMETERS align, norm, broad -----

0.025 0.004 1.500 3.000 !norm,+/-,Exp.shift,+/-  
Y 45000 ! Broad? Tot trials  
-3.500 2.000 ! E.Fermi -- +/-  
1.600 0.400 ! Broad -- +/-  
0.800 0.600 ! Broad\_Exp -- +/-  
Y ! mf?  
1 ! # plasmons  
14.000 10.000 10.000 8.000 ! En. +/- , Ampl. +/-  
N 0.02 70.0 0.02 !Extra Exc? ampl,energy,width

----- OTHER OPTIONS -----

Y 6 ! Correlation links? How many?  
1 7  
4 10  
2 8  
5 11  
16 22  
19 25 ! between which params ?  
Y N !Slope?,Concerted motions?  
N 30.000 10.000 !weight?(N,A,T,G) xcentr, width

SET PRINT 3

SEEK 90

!SCAN 13 10 -.2 .2

!FIX 2.

SIMPLEX 100 0.03

CAIJ.FCN -2

**Control file** – general information – methods for fitting – It is always called **COMMAND.MIN**

We get several outputs, the most important is the **BEST** file where you find all information about the fit

**FPMS**

- Input file must be named as "data.ms".
- In principle the way to use parameters are the same as "data.ms" of program POTGEN or "DATA.\*\*\*\*" of MXAN.
- For non-MT calculation, you should add,  
NMT = "NM"  
truncate = .true. (voronoi)  
or  
truncate = .false. (ASA)