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## **Experimental powder pattern**

A powder diffraction pattern can be recorded in numerical form for a discrete set of scattering angles, times of flight or energies. We will refer to this scattering variable as : T. The experimental powder diffraction pattern is usually given as three arrays :

$$\left\{T_i, y_i, \sigma_i\right\}_{i=1,2,\dots,n}$$

The profile can be modelled using the calculated counts:  $y_{ci}$ at the *i*th step by summing the contribution from neighbouring Bragg reflections plus the background.

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$$y_{ci} = \sum_{\{\mathbf{h}\}} I_{\mathbf{h}} \Omega(T_i - T_{\mathbf{h}}) + b_i$$

The symbol {h} means that the sum is extended only to those reflections contributing to the channel ``i" .

This should be taken into account (resolution function of the diffractometer and sample broadening) before doing the actual calculation of the profile intensity.

This is the reason why some Rietveld programs are run in two steps (C) % = ]]»

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to the diffraction pattern

Several phases ( $\phi = 1, n_{\phi}$ ) contributing

$$y_{ci} = \sum_{\phi} s_{\phi} \sum_{\{\phi \mathbf{h}\}} I_{\phi, \mathbf{h}} \ \Omega(T_i - T_{\phi, \mathbf{h}}) + b_i$$

Several phases ( $\phi = 1, n_{\phi}$ ) contributing to several  $(p=1,n_p)$  diffraction patterns

$$y_{ci}^{p} = \sum_{\phi} S_{\phi}^{p} \sum_{\{\phi\mathbf{h}\}} I_{\phi,\mathbf{h}}^{p} \Omega^{p} (T_{i} - T_{\phi,\mathbf{h}}) + b_{i}^{p}$$

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$$y_{ci} = \sum_{\{\mathbf{h}\}} I_{\mathbf{h}} \Omega(T_i - T_{\mathbf{h}}) + b_i$$
$$I_{\mathbf{h}} = S \left\{ L p O A C F^2 \right\}_{\mathbf{h}}$$

Integrated intensities are proportional to the square of the structure factor F. The factors are: Scale Factor (S), Lorentz-polarization (Lp), preferred orientation (O), absorption (A), other "corrections" (C)

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**The Structure Factor contains** the structural parameters (isotropic case)  $F(\mathbf{h}) = \sum_{j=1}^{n} O_{j} f_{j}(\mathbf{h}) T_{j} \sum_{s} exp\left\{2\pi i \left[\mathbf{h}\left\{S \mid \mathbf{t}\right\}_{s} \mathbf{r}_{j}\right]\right\}$  $\mathbf{r}_{j} = (x_{j}, y_{j}, z_{j})$  (j = 1, 2, ...n) $T_j = \exp(-B_j \frac{\sin^2 \theta}{\lambda^2})$ 

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The Structure Factor in  
complex cases
$$F(\mathbf{h}) = \sum_{j=1}^{n} O_j f_j(h) T_j \sum_{s} g_j(\mathbf{h}_s) exp \left\{ 2\pi i \left[ \mathbf{h} \left\{ S | \mathbf{t} \right\}_s \mathbf{r}_j \right] \right\}$$
 $\mathbf{h}_s = \begin{pmatrix} h \\ k \\ l \end{pmatrix}_s = S_s^T \begin{pmatrix} h \\ k \\ l \end{pmatrix}$  ( $s = 1, 2, ...N_G$ ) $g_j(\mathbf{h}_s)$ Complex form factor of object j  
Anisotropic DPs  
Anharmonic DPs

The peak shape function of powder diffraction patterns contains the Profile Parameters

$$\Omega(x_{\mathbf{h}i},\beta_{\mathbf{P}}) = \Omega(T_i - T_{\mathbf{h}},\beta_{\mathbf{P}})$$

$$\Omega(x) = g(x) \otimes f(x) = instrumental \otimes intrinsic profile$$

$$\int_{-\infty}^{+\infty} \Omega(x) dx = 1$$

In most cases the observed peak shape is approximated by a linear combination of Voigt (or pseudo-Voigt) functions

$$\Omega(x) \approx \sum L(x) \otimes G(x) = \sum V(x)$$

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#### **Properties of the Voigt function**

 $V(x) = V_1(x) \otimes V_2(x)$ 

The Voigt function has proven to be a very good experimental approximation in many cases

$$\beta_{L} = \beta_{1L} + \beta_{2L}$$
Lorentzian breadths simply  
have to be summed  

$$\beta_{G}^{2} = \beta_{1G}^{2} + \beta_{2G}^{2}$$
Gaussian breadths have to be  
summed quadratically  

$$\beta_{fL} = \beta_{hL} - \beta_{gL}$$
Correction for  
instrumental broadening  

$$\beta_{fG}^{2} = \beta_{hG}^{2} - \beta_{gG}^{2}$$
Correction for  
instrumental broadening

The Rietveld Method consist of refining a crystal (and/or magnetic) structure by minimising the weighted squared difference between the observed and the calculated pattern against the parameter vector:  $\beta$ 

$$\chi^{2} = \sum_{i=1}^{n} w_{i} \left\{ y_{i} - y_{ci}(\beta) \right\}^{2}$$

$$W_{i} = \frac{1}{\sigma_{i}^{2}}$$

$$\sigma_{i}^{2}$$
 is the variance of the "observation"  $y_{i}$ 

However, the Rietveld Method can be easily extended by using, instead of the traditional  $\chi^2$ (least squares), another Cost Function to be minimised against the parameter vector  $\beta$ 

$$Cost = \sum_{i=1}^{n} F(\{y_i - y_{ci}(\beta)\})$$

$$Cost = -\log(Likelihood)$$

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## Least squares: Gauss-Newton (1)

Minimum necessary condition:  $\frac{\partial \chi^2}{\partial \beta} = 0$ 

A Taylor expansion of  $y_{ic}(\beta)$  around  $\beta_0$  allows the application of an iterative process. The shifts to be applied to the parameters at each cycle for improving  $\chi^2$  are obtained by solving a linear system of equations (normal equations)

$$\mathbf{A}\boldsymbol{\delta}_{\boldsymbol{\beta}_{0}} = \mathbf{b}$$

$$A_{kl} = \sum_{i} W_{i} \frac{\partial y_{ic}(\boldsymbol{\beta}_{0})}{\partial \beta_{k}} \frac{\partial y_{ic}(\boldsymbol{\beta}_{0})}{\partial \beta_{l}}$$

$$b_{k} = \sum_{i} W_{i}(y_{i} - y_{ic}) \frac{\partial y_{ic}(\boldsymbol{\beta}_{0})}{\partial \beta_{k}}$$

#### Least squares: Gauss-Newton (2)

The shifts of the parameters obtained by solving the normal equations are added to the starting parameters giving rise to a new set

$$\boldsymbol{\beta}_1 = \boldsymbol{\beta}_0 + \boldsymbol{\delta}_{\boldsymbol{\beta}_0}$$

The new parameters are considered as the starting ones in the next cycle and the process is repeated until a convergence criterion is satisfied. The variances of the adjusted parameters are calculated by the expression:

$$\sigma^{2}(\beta_{k}) = (\mathbf{A}^{-1})_{kk} \chi_{\nu}^{2}$$
$$\chi_{\nu}^{2} = \frac{\chi^{2}}{N - P + C}$$

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# Least squares: a local optimisation method

• The least squares procedure provides (when it converges) the value of the parameters constituting the local minimum closest to the starting point

• A set of good starting values for all parameters is needed

• If the initial model is bad for some reasons the LSQ procedure will not converge, it may diverge.



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#### Needs for precise refinements and microstructural effects

Precise refinements can be done with confidence only if the intrinsic and instrumental peak shapes are properly approximated.

At present

 $\Rightarrow$  The approximation of the intrinsic profile is mostly based in the Voigt (or pseudo-Voigt) function

 $\Rightarrow$  The approximation of the instrumental profile is also based in the Voigt function for constant wavelength instruments

⇒ For TOF the instrumental+intrinsic profile is approximated by the convolution of a Voigt function with back-to-back exponentials or with the Ikeda-**Carpenter function.** 

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#### Example: General $2\theta$ dependence of the instrumental broadening (determined by a standard sample)

$$H_{hG}^{2} = (U_{f} + (1 - \xi_{f})^{2} D_{fST}^{2}(\boldsymbol{a}_{D})) \tan^{2} \theta + \frac{I_{fG}}{\cos^{2} \theta} + H_{gG}^{2}$$
$$H_{hL} = (X_{f} + \xi_{f} D_{fST}(\boldsymbol{a}_{D})) \tan \theta + \frac{[Y_{f} + F_{f}(\boldsymbol{a}_{S})]}{\cos \theta} + H_{gL}$$

The Gaussian and Lorentzian components of the instrumental Voigt function are interpolated between empirically determined values. If needed, axial divergence is convoluted numerically with the resulting profile.

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#### Microstructural effects and peak shapes (Rietveld)

**Recent developments:** 

⇒ Anisotropic peak broadening (strain/size effects): quartic forms in hkl (dislocations, micro-twinning composition fluctuations) spherical harmonics (complex size/microstrain effects) special reflections (stacking faults, antiphase domains, polytypes)

 $\Rightarrow$  *hkl* - dependent shifts with respect to Bragg positions: special reflections, quartic forms, ...

⇒ New sample profiles: Linear combination of pseudo-Voigt functions to mimic log-normal and gamma size distributions (Popa et al. J.Appl. Cryst 35, 2002, 338-346)

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### Problems when modeling the peak shape, a real case: low resolution neutron powder diffractometers

#### D20 at ILL:

A diffraction pattern can be collected in less than a second!

Large graphite monochromator with a quite low take-off angle  $(2\theta_{\rm M}\approx 40^{\circ})$ 

This implies that at high angle the peaks are broad and have a strange peak shape (that can be reproduced quite precisely by ray tracing or Monte Carlo simulations of the instrument !)

Test Na2Ca3Al2F14 D20 (ncad20\_24\_10\_pr7/NCF24\_10)

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20000

16000

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The inability to model peak shapes properly

introduces a "systematic error" in the data treatment affecting the structural parameters and the estimation of their uncertainties

#### How can we solve this problem?

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#### Future developments for Rietveld analysis: the treatment of the peak shape

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#### Further step on complexity:

⇒ Fundamental parameters approach

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- $\Rightarrow$  Numerical instrumental profile (when needed)
- ⇒ Local convolution with analytical sample profile using FFT or interpolated direct convolution

$$\Omega(x) = FT^{-1}[G(t)F(t)] = g(x) \otimes f(x)$$

This is partially performed (with analytical functions) in the CCSL based code at ISIS and in FullProf/GSAS for the TOF case:  $V(x) \otimes IK(x)$ In TOPAS the fundamental parameters approach is fully implemented

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# Future developments for Rietveld analysis: the treatment of the peak shape

⇒ Different components of both instrumental and sample profile functions are just multiplied in the Fourier space. ⇒ The global G(t) may be provided in the instrumental resolution file in different forms, depending if it can be approximated by analytical functions or not.

$$G(t) = G_1(t)G_2(t)G_3(t)...$$
  

$$F(t) = F_1(t)F_2(t)F_3(t)...$$

This procedure is faster than the direct convolution using numerical integration when the number of points per profile is greater than  $\sim 64$ .

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# The Rietveld algorithm: (do over points/reflections)

```
Subroutine calc_patterns_TOF(n_pat)
.....
DO i=1,npts(n_pat)
ini=code_contribution(i,n_pat,"ini")
fin=code_contribution(i,n_pat,"fin")
IF(iprev <= fin) THEN
DO j=iprev,fin
CALL calcul_tof(j,n_pat)
END DO
iprev=MAX(iprev,fin+1)
END IF
CALL summat_tof(i,n_pat,ini,fin)
END DO
return
End Subroutine calc_patterns_TOF
```

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# Do for N\_cycles ← may be done in parallel Do for Patterns ← may be done in parallel Do for contributing reflections

Do for contributing reflections Calculate broadening w.r.t to IRF Calculate structure factors+derivatives Sum contributions (LSQ matrix + vector) calculate profile for current point and reflections contributing to it (convolution ⇒ neighbours needed) profile derivatives End do reflections End do points in Patterns End do Patterns Invert LSQ matrix and update the free parameters Tests for convergence (if convergence is reached exit!) End do N\_cycles With 2007

# May the Rietveld algorithm be improved?

New ideas are needed to improve the efficiency:

New data structures? Store individual peak shapes? Change the order of loops? Modularise different parts of the calculations?



With the forthcoming high performance instruments, and increasing complexity, we need an improvement of the algorithms for handling Rietveld refinements if we want to preserve interactivity. Options:

- $\Rightarrow$  Develop small specialized Rietveld programs
- ⇒ Maintain the possibility of general treatment and improve the efficiency by making strong changes on the Rietveld codes.
- $\Rightarrow$  A combination of both ...

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## **Conclusions and perspectives**

The increasing complexity of instruments and the necessity of better refinements call for collaborative teams for improving the existing software and develop new tools.

A list of tasks and priorities is needed to undertake a rational software development.

This opens new opportunities for young people wishing to dedicate their scientific activities to Crystallographic Computing.



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