IUCr Congress 2017. XAFS Tutorial for Crystallographers. Hyderabad Monday 21st August XAFS Theory: XANES and EXAFS Spectra.

### Chris. Chantler Prof., FAIP, Chair CXAFS University of Melbourne Victoria 3010, Australia chantler@unimelb.edu.au



http://optics.ph.unimelb.edu.au/~chantler/ home.html

MELBOURNE

SESSION 1. Chair: Chris Chantler (Australia)

9:00 X-ray absorption spectroscopy for beginners. Overview - Farideh Jalilevand (Canada)

9:30 XAFS Theory: XANES and EXAFS Spectra – Chris Chantler (Australia)

10:00 Coffee

**SESSION 2.Chair: Farideh Jalilevand (Canada)** 

10:30 XAS Experiment - Steve Heald (USA)

**11:15** EXAFS Analysis - Boopesh Mishra (UK)

12:00 Summary of Software and Panel Discussion - Chris Chantler (Australia) and panel

12:20 Applications I: Materials - Guiliana Aquilanti (Elettra, Italy)

13:05 lunch

**SESSION 3.Chair: Christopher Chantler (Australia)** 

14:55 Applications II: Biology - Ritimukta Sarangi (Stanford)

15:40 Applications III: Industry - Sohini Basu (Honeywell UOP, India)

**16:00** Applications III: Environmental - Hugh Harris (Australia)

16:45 Proposal Preparation for Synchrotron Beamtime and Closing Notes – Chris Chantler (Australia) [Bruce Ravel]

**16:55 Workshop Close. Congress Reception!** 

Thanks for this initiative: **IUCr Commission on XAFS** International X-ray Absorption Society **IUCr Local Organising Committee IUCr Congress Farideh Jalilehvand, Bruce Ravel, Dibyendu Bhattacharyya** All organisers, executive, etc. Session leaders and presenters at IUCr Congress **Continuing engagement, collegial development** IUCr 2017. XAFS Tutorial. C.T. Chantler





### International Tables for Crystallography Volume I X-ray Absorption Spectroscopy



Acta Crys

Editors: Chris Chantler, Federico Boscherini, Bruce Bunker



Acta Cryst

Acta Cryst

- 1. Introduction
- 2. XAS Theory
- 3. Experimental Methods
- 4. Pre-processing of Experimental Data
- 5. Overview of Analysis of Experimental Data
- 6. Data Reduction, Theoretical Prediction and Data Analysis
- 7. Exchange of Data and Processed Outputs, Deposition8. Applications









## Volume I

- What [old and new] things can we do with XAS that we could not before, how can we do them, and what standards can we use?
- Research in core physics and condensed matter science are increasingly relevant for diverse fields
- Applications in chemistry, engineering and biological sciences, linking to experimental research at synchrotrons, reactors and specialized facilities
- A plethora of different approaches are popular in the literature
- The Volume will hope to capture great achievements and value by representation of leading groups from Europe, America, Asia, Australia plus elsewhere!
- traditional and new experiments, theory, analysis

### **Recent New Fields**

- Resonant Inelastic X-ray Scattering (Glatzel, numerous other groups, several important variant techniques) - scattering and resonant analysis; huge information content; complex but successful theory; insight well below hole width limits
- Ultra-fast XAS (Frahm, XFELs, MaxLab; SLS; Tsukuba) dynamic processes; direct observation of catalysis; phase changes; ...
- X-ray Emission Spectroscopy electronic structure, X-ray Raman spectroscopy...
- X-ray Excited Optical Luminescence site specificity, optical properties
- High Energy Resolution Fluorescence reduced lifetime broadening





 Why XAFS?
 What is XAFS?
 How does XAFS work?
 Difficulties
 Links with Crystallography Realisation: A: Absorption Realisation: B:Fluorescence

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 $k (Å^{-1})$ 

### 1. Why XAFS?

- Crystallography periodicity, symmetry and mean structure of perfect crystals [X-ray; Electron Diffraction; Neutron Diffraction]
- Nanocrystallites in advanced synchrotron beams; or [Xray] Powder Diffraction
- TEM etc. of surfaces or slices [destructive].
- Great difficulties for disordered systems; solutions; dilute systems; local order; dynamic bond lengths; active centres.
- X-ray Absorption Fine Structure [XAFS] deals directly with these questions and more.
- Complementary for complex systems, organometallics, bioactive systems, ideal crystals or metals

### 1. Why XAFS?

- Recent advances in synchrotron techniques,
- new opportunities in organometallic materials,
- dynamic bond investigation,
- thermal dependence of structure [especially disordered],
- complex phase systems
- cluster studies
- >12000 papers in the last 5 years.





Application to physics, earth science, chemistry and biology Ascone et al. (2009). Journal of Synchrotron Radiation 16, 413-42

### **Spectroscopy**





### Changes in electron state when a photon is absorbed by an atom.

These mechanisms enable: a) x-ray photoelectron spectroscopy b) x-ray absorption spectroscopy c) x-ray emission spectroscopy d) auger electron spectroscopy

### X-ray and photon interaction processes





Interactions include:

photoelectric absorption  $\sigma_{pe}$ ,  $\tau$ 

inelastic scattering (Compton Effect) σ<sub>inel</sub> (incoh)

including X-ray fluorescence

elastic scattering (Rayleigh)  $\sigma_{el}$  (coh)

pair production  $\sigma_{pp}(\kappa)$ 

### Links: X-ray Absorption Fine Structure and Crystallography Photoelectric Absorption; Coherent (elastic) Scattering; Inelastic (Compton) Scattering [e.g. Hubbell...]

<u>Cross-sections (barns/atom)</u>:  $\sigma_{tot} = \tau_{PE} + \sigma_{coh} + \sigma_{incoh} + \kappa_{h} + \kappa_{e} + \sigma_{p.n.}$  $\tau_{PF}(E) = 2hcr_{a} \operatorname{Im}(f)/E$ <u>X-ray Absorption</u>: mass absorption coefficient (cm<sup>2</sup>/g):  $[\mu_{PE} / \rho](E) = \tau_{PE} / (uA)$ linear absorption coefficient (cm<sup>-1</sup>):  $\mu_{PE} = [\mu_{PE} / \rho]\rho$ Coherent (Rayleigh) scattering:  $\sigma_{coh,Rayleigh} = \pi r_e^2 \int (1 + \cos^2 \theta) f^2(q,Z) d(\cos \theta)$ Coherent (Bragg-Laue) scattering:  $\sigma_{coh,BL} = \left(\frac{r_e^2 \lambda^2}{2NV_e}\right) \sum_{\mu}^{-1} \left[ \left(\frac{1 + \cos^2 \theta}{2}\right) m d |F|^2 \right]$  $F(hkl) = \sum f_{i} e^{-M_{j}} e^{2\pi i(hx_{j} + ky_{j} + lz_{j})}, (TDS = 0)$ Coherent (Thermal Diffuse) Scattering:  $\sigma_{coh,TDS} = \left(\frac{r_e^2 \lambda^2}{2NV_c}\right) \sum_{H} \left[ \left(\frac{1+\cos^2\theta}{2}\right) m d|F|^2 \left\{1-e^{-2M}\right\} \right]_{H} \right]_{H}$ Inelastic (Compton):  $\sigma_{incoh,Compton} = \pi r_e^2 \int_{-1}^{1} \left(\frac{1+\cos^2\theta + \frac{k^2(1-\cos\theta)^2}{1+k(1-\cos\theta)}}{(1+k(1-\cos\theta))^2}\right) S(q,Z) d(\cos\theta), k = \frac{\hbar\omega}{mc^2}$  $I_{incoh} = I_e \left( \sum_{j=1}^{Z} (1 - f_j^2 - \sum_{k \neq j} \int \psi_j^* \psi_k e^{i\mathbf{q}\cdot\mathbf{r}} d\mathbf{r}) \right) = I_e S(q, Z)^{\left( \frac{1 + \cos^2 \theta}{2} \right)}$ Klein-Nishina (free electron)  $I_e = I_0 r_e^2 \left[ \frac{1 + \cos^2 \theta}{2} \right]$  $\kappa_n$  nuclear-field pair production  $\kappa_e$  triplet production  $\sigma_{p.n.}$  photonuclear absorption

### Links: X-ray Absorption Fine Structure and Crystallography ATTENUATION & THE (X-RAY) ATOMIC FORM FACTOR: One Example of usage: FFAST:

$$\mu_{PE}(E) = f_2(E) 2hcr_0 / E$$

$$[\mu_{PE} / \rho](E) = \sigma_{PE} / (uA)$$

Im(f) = f''(E) = f\_2(E) = 
$$\frac{E\mu_{PE}(E)}{2hcr_e}$$

$$\operatorname{Re}(f) = f_0 + f' + f_{NT}, f' = f_1 + f_{rel} - Z$$
$$f_0(q, Z) = 4\pi \int_0^\infty \frac{\rho(r) \sin(qr) r^2 dr}{qr}$$

http://physics.nist.gov/PhysRefData/FFast /Text/cover.html

Chantler, CT, Olsen, K, et al. (2005) X-Ray Form Factor, Attenuation & Scattering Tables (v2.1) [Online]; Chantler, CT, JPhysChemRefData 29(4), 597-1048 (2000); Chantler, CT, JPhysChemRefData 24, 71-643 (1995).



Links: X-ray Absorption Fine Structure and Crystallography (X-RAY) ATOMIC FORM FACTOR: Resonant scattering amplitude of X-rays by charge (electron) density

$$\operatorname{Re}(f) = f_0 + f' + f_{NT}, f' = f_1 + f_{rel} - Z$$
$$f_0(q, Z) = 4\pi \int_0^\infty \frac{\rho(r) \sin(qr) r^2 dr}{qr}$$
$$f'(E, Z) = f'(\infty) - \frac{2}{\pi} P \int_0^\infty \frac{\varepsilon' f''(\varepsilon')}{E^2 - (\varepsilon')^2} d\varepsilon'$$

Im(f) = f''(E) = 
$$f_2(E) = \frac{E\mu_{PE}(E)}{2hcr_e}$$

'normal' coherent scattering factor

'anomalous' scattering factor

<u>Utility</u>: X-ray Diffraction experiments, Crystallography...

$$F(hkl) = \sum_{j} f_{j} e^{-M_{j}} e^{2\pi i (hx_{j} + ky_{j} + lz_{j})}, (TDS = 0)$$

Electronic wavefunction distribution, bonding...

$$\Delta \rho(x, y, z) = \frac{1}{V} \sum_{k} \sum_{k} \sum_{l} \Delta F(hkl) e^{-2\pi i (hx + ky + lz)}$$

VUV research, multilayer modelling, critical angle spectroscopy...  $n_r = n + ik = \sqrt{\varepsilon} = 1 - \delta - i\beta = 1 - \frac{r_0}{2\pi} \lambda^2 \sum_j n_j f_j$   $n_j$  atom number density;  $r_0$  classical electron radius Transmission, attenuation experiments...  $\mu_{PE}(E) = f_2(E)2hcr_0 / E$ Electron form factor (Mott-Bethe)...  $f^B(q,Z) = \frac{me^2}{2\pi\hbar^2\varepsilon_0} \left\{ \frac{Z - f(q,Z)}{q^2} \right\}$ 

## XAFS Theory: XANES and EXAFS Spectra XAFS, XANES, XERT





### 2. What is X-ray Absorption Fine Structure XAFS?

- sharp oscillations in Absorption above a absorption edge [K, LI, LII, LIII, M...] of an element [Fe, Cu, Ni, C, Mo, Au, ...] in a material
- creation of electron holes in 1s, 2s, 2p1/2, 2p3/2
- The material may be an ideal crystal or metal, a nanocrystal or powder, or a non-ideal mixture or dilute solution.

See IUCr Commission on XAFS Definitions: <u>http://</u> <u>www.iucr.org/resources/commissions/xafs/xafs-related-</u> <u>definitions-for-the-iucr-dictionary</u> See also: <u>http://</u> <u>www.iucr.org</u>/

### 2. What is XAFS?

- modulation of absorption coefficient due to chemical state & structure of immediate surroundings.
- •'near edge' region (XANES or NEXAFS) to ~50 eV above the absorption edge
- 'extended' region (EXAFS) giving oscillations in the absorption coefficient from ~50 eV.
- •XANES (X-ray Absorption Near Edge Spectroscopy) for X-ray edges (~1 keV and above); NEXAFS (Near-edge Xray Absorption Fine Structure) for soft X-ray edges.
- •Spectral features before the main absorption edge - 'pre-edge' features are associated with transitions to bound states.



## MODERN QUANTITATIVE THEORY – FMS MULTIPLE SCATTERING

Start with Fermi's Golden Rule:

$$\mu \propto \sum_{f} \left| \left\langle \boldsymbol{\psi}_{i} \right| d \left| \boldsymbol{\psi}_{f} \right\rangle \right|^{2} \delta \left( E_{f} - E_{i} - \hbar \boldsymbol{\omega} \right)$$

- Final (excited) state wavefunctions are difficult to treat directly in an arbitrary solid
- Use a trick: expand out...

$$\boldsymbol{\mu} \propto \langle \boldsymbol{\psi}_i | d \sum_f | \boldsymbol{\psi}_f \rangle \langle \boldsymbol{\psi}_f | \boldsymbol{\delta} (E_f - E_i - \hbar \boldsymbol{\omega}) d^* | \boldsymbol{\psi}_i \rangle$$

$$\mu \propto -\frac{1}{\pi} \operatorname{Im} \langle \Psi_i | dG(\omega) d^* | \Psi_i \rangle$$

## MODERN QUANTITATIVE THEORY - FMS

Green's function propagator:

 $G = [E - H_0 - V]^{-1}$ 

$$= [E - H_0]^{-1} \left\{ 1 + V(E - H_0)^{-1} + [V(E - H_0)^{-1}]^2 + \dots \right\}$$

 $G = G_0 + G_0 V G_0 + G_0 V G_0 V G_0 + \dots$ 

 Muffin-tin approximation – V's become site based block diagonal matrices T

$$G = G_0 + G_0 T G_0 + G_0 T G_0 T G_0 + \dots$$

 $G(\omega) = [1 - G_0 T]^{-1} G_0$ 

### FULL MULTIPLE SCATTERING

\* c/o Josh Kas, University of Washington



## FDMNES - METHOD OF CALCULATION

### •Sum over transition amplitudes

$$M_{gf} = \left\langle \Psi_f \right| \varepsilon \cdot r \left( 1 + \frac{i}{2} k \cdot r \right) \left| \Psi_g \right\rangle$$

To evaluate optical cross section

$$\sigma = 4\pi^2 \alpha \hbar \omega \sum_{f,g} \left| M_{gf} \right|^2 \delta \left( \hbar \omega - E_f + E_g \right)$$

Require wavefunctions for final and initial states
Initial state is a ground state – calculated atomistically with relativistic corrections
Final state determined using the FDM

## FDMNES - METHOD OF CALCULATION

•Cluster broken up into three zones – (1) near cores (small), (2) interstitial, (3) outside cluster

•Wave functions in zone 1 constructed from spherical harmonics, in zone 3 from Neumann and Bessel functions



 In zone 2 the FDM is used to determine the final-state photoelectron wave function

## FDMNES - METHOD OF CALCULATION

 Potential in interstitial region may be determined from DFT internally or externally (WIEN2k), or from electron density values with exchange+correlation correction

•Once potential is found, evaluate wavefunctions using discretized Schrodinger Equation:

$$\nabla^2 \boldsymbol{\psi}_i = \frac{1}{h^2} \left( \frac{4}{3} \sum_{j,\varepsilon} \boldsymbol{\psi}_j^{\varepsilon} - \frac{1}{12} \sum_{j,\varepsilon} \boldsymbol{\psi}_j^{\varepsilon\varepsilon} - \frac{15}{2} \boldsymbol{\psi}_i \right)$$
$$\left( -\nabla^2 + V_i - E \right) \boldsymbol{\psi}_i + \sum_j \left( -\nabla^2 \boldsymbol{\psi}_j \right) = 0$$

## FDMNES - EXAMPLES AND HANDS-ON

### /fdmnes/fdmfile.txt

#### fdmfile.txt ! General indata file for FDMNES ! with indata files examples 1 entree/test\_stand/cu\_inp2.txt entree/test\_stand/vo6 inp.txt entree/test\_stand/vo6 conv inp.txt entree/test\_stand/vo6\_nodipole\_inp.txt entree/test\_stand/feo6\_inp.txt entree/test stand/ni inp.txt entree/test\_stand/ni\_mg\_inp.txt entree/test\_stand/v2o3\_inp.txt entree/test\_stand/gan\_inp.txt entree/test\_stand/fe3o4\_inp.txt entree/test\_stand/fe3o4 dd inp.txt entree/test\_stand/cr\_inp.txt entree/test\_stand/cr\_conv\_inp.txt entree/test\_stand/ba2znuo6 inp.txt entree/test\_stand/ca3co2o6\_inp.txt entree/test\_stand/fe2o3\_inp.txt Filout entree/test\_stand/fe2o3 selec inp.txt entree/test\_stand/fe2o3\_scf\_inp.txt entree/test stand/fe2o3 hub inp.txt entree/test\_stand/fe bio inp.txt entree/test stand/ndmg inp.txt entree/test\_stand/pt13 inp.txt entree/test\_stand/mult\_inp.txt entree/test stand/cuo inp.txt

### •Run FDMNES.exe

## /fdmnes/entree/ cu inp.txt

Fdmnes indata file

Calculation for the copper K-edge in copper cfc

! Finite difference method calculation with convolution

xanout/test\_stand/cu

Range -10. 0.2 0. 0.5 10. 1. 4	<pre>! Energy range of calculation (eV) 0. ! first energy, step, intermediary energy, step, last energy</pre>
Radius 3.0	! Radius of the cluster where final state calculation is performed ! For a good calculation, this radius must be increased up to 6 or 7 <u>Angstroems</u>
Crystal 3.61 3.61 3.61 90. 29 0.0 0.0 0.0 29 0.5 0.5 0.0 29 0.5 0.0 0.5 29 0.0 0.5 0.5	<pre>! Periodic material description (unit cell) 90. 90. ! a, b, c, (Angstroem) alpha, beta, gamma (degree)</pre>

cu inp.txt

! Convolution keyword : broadening with a width increasing versus energy as an arctangent

Convolution

End

### FDMX OVERVIEW AND PURPOSE



### 3. How does XAFS work?

### Local Structure Guess



Metallic Molybdenum: BCC Crystal

### Model of XAFS Spectrum (XAFS Equation)

**Numbers** 

Program?	$\chi_{th}(k) = \sum_{j} r_{j}$ $r_{j} = (1 + \alpha)r_{0}$	$\chi_{th}(k) = \sum_{j} N_{j} S_{0}^{2} F_{j}(k) \frac{\sin(2kr_{j} + \phi_{j}(k))}{kr_{j}^{2}} e^{-2\sigma_{j}^{2}k^{2}} e^{-2r_{j}/\lambda_{j}(k)}$ $r_{j} = (1 + \alpha)r_{0,j}$					
	Expansion Coefficient	Bac Am	<b>kscatterin</b> <b>plitude an</b> Many Body	g d Phas	Mean <sup>e</sup> Free Path		
	Distance To Coordination Shell		Factor Coordinat	Debye Factor	Debye-Waller Factor on		

Inclusion of thermal disorder

$$\chi_j(k) = \dots e^{\left(-2\sigma_j^2 k^2\right)}$$

•Require *relative* mean-square displacement of neighbour atoms along a photoelectron scattering path  $R_i$ 

$$\sigma_j^2 = \left\langle (u_0 \cdot R_j)^2 \right\rangle + \left\langle (u_j \cdot R_j)^2 \right\rangle - 2 \left\langle (u_0 \cdot R_j) \cdot (u_j \cdot R_j)^2 \right\rangle$$

•For *N* atoms of mass *M* in a crystal with resonant phonons of energy  $\hbar \omega_{\alpha\lambda}$ , express as:

$$\sigma_j^2 = \frac{\hbar}{NM} \sum_{q\lambda} (e_{q\lambda} \cdot R_j)^2 \frac{1}{\omega_{q\lambda}} \left( \coth\left(\frac{\omega_{q\lambda}}{2k_BT}\right) \right) \left[ 1 - \cos(q \cdot R_j) \right]$$

\* Gruneisen & Goens, *Z. Phys.* 29 141(1924)
\* Kittel, Introduction to Solid State Physics, Wiley, 1996

**FDMX METHOD OF CALCULATION** •Substituting  $\omega_{q\lambda}$  values from Debye Model, and Taylor expanding coth term enables "simple" equation for a correlated Debye model: Beni-Platzmann theory

$$\sigma_{j}^{2} = \frac{6\hbar}{m\omega_{D}} \left[ \frac{1}{4} + \left(\frac{T}{\theta_{D}}\right)^{2} D_{1} \right] - \frac{6\hbar}{m\omega_{D}} \left\{ \frac{1 - \cos(q_{D}r_{j})}{2(q_{D}r_{j})^{2}} + \left(\frac{T}{\theta_{D}}\right)^{2} \left[ D_{1} - \frac{1}{3!} \left(q_{D}r_{j}\frac{T}{\theta_{D}}\right)^{2} D_{3} + \frac{1}{5!} \left(q_{D}r_{j}\frac{T}{\theta_{D}}\right)^{4} D_{5} - \dots \right] \right\}$$

•Enables efficient calculation of Debye-Waller factor  $\sigma^2$  for each atom in a cluster

•Only approximates density of phonon states – full account may be implemented with DFT but much more complex/inefficient!

•First order correction from nearest-neighbour (dominant) Debye-Waller factor  $\chi_i(k) \rightarrow \chi_i(k) e^{(-2\sigma_j^2k^2)}$ 

Accurate at high energies – area of maximum importance

 Underestimate at low energies, but not a dominant effect in this region (i.e. compare with IMFPs, potential shape etc.)



 Inelastic scattering causes an energy-dependent broadening of the XAFS

$$\Gamma(E) = \frac{\hbar}{\lambda(E)} \sqrt{\frac{2E}{m}} + \Gamma_{H}$$



Core-hole lifetime also included
Strong effect beyond 10 eV

•Use many-pole optical data model and generalized theory of Tanuma et al.

\* S. Tanuma et al. *Surf. Interface Anal.* **43** 689 (2011)

# •IMFP determined from electron energy loss function (ELF)

$$\lambda^{-1}(E) = \frac{\hbar}{a_0 \pi E} \int_{0}^{\frac{E-E_F}{\hbar}} \int_{q_-}^{q_+} \frac{1}{q} \operatorname{Im}\left[\frac{-1}{\varepsilon(q,\omega)}\right] dqd\omega$$
$$q \pm = \sqrt{\frac{2mE}{\hbar^2}} \pm \sqrt{\frac{2m}{\hbar^2}(E - \hbar\omega)}$$

### •Electron ELF built from transform of optical dielectric function $\varepsilon(0,\omega)$

$$\operatorname{Im}\left[\frac{-1}{\varepsilon(0,\omega)}\right] = \sum_{i} A_{i} \operatorname{Im}\left[\frac{-1}{\varepsilon_{FEG}(0,\omega;\omega_{p}=\omega_{i})}\right]$$



## FDMX - EXAMPLES AND HANDS-ON

# •Further examples of FDMX:

Molybdenum:





Ferrocene:

# X-ray Absorption Spectroscopy how to measure

### XANES, EXAFS, ...

Most fundamental technique is a transmission mode



 $\mu t (E) = ln (i_o / i)$ 

You measure attenuated beam intensity, that" exponentially" decreases

 $\mu t(E) = F / i_0$ 

You measure emitted beam intensity Which "linearly" proportional to conc.



The fluorescence intensity  $I_f(E)$  accepted by a detector with a solid angle of  $\Omega/4\pi$  is:

$$I_{f} = \frac{\frac{fI_{0}\Omega[\mu/\rho]_{pe}^{*}}{4\pi\cos\theta_{inc}}}{\frac{[\mu/\rho]}{\cos\theta_{inc}} + \frac{[\mu_{f}/\rho]}{\cos\theta_{out}}} \left( A - \exp\left(-\frac{\left[\frac{\mu}{\rho}\right]\left[\rho t\right]}{\cos\theta_{inc}} - \frac{\left[\frac{\mu_{f}}{\rho}\right]\left[\rho t\right]}{\cos\theta_{out}}\right) \right)$$
$$I_{f}(E) \approx I_{0} \times 8.8 \times 10^{-4} \times 0.012 \approx 1 \times 10^{-5} I_{0}, \text{ lower than } I_{0} \text{ by 5 ordels.}$$

## EXPERIMENTAL ANALYSIS AND SYSTEMATICS

- Beer Law:  $I_1 = I_0 e^{(-\mu t)}$
- Dark current (no beam through each detector)

$$\left[\frac{\mu}{\rho}\right] \left[\rho t\right] = -\ln\left(\frac{I_1 - D_1}{I_0 - D_0}\right)$$

• Air path offset – blank in sample stage

$$\left[\frac{\mu}{\rho}\right] \left[\rho t\right]_{s} = \left[\frac{\mu}{\rho}\right] \left[\rho t\right]_{s+a} - \left[\frac{\mu}{\rho}\right] \left[\rho t\right]_{a}$$









What are the uncertainties? How much information content remains? How much has been lost for hypothesis testing?





## **Software Packages**

3. How does XAFS work?

$$\chi_{th}(k) = \sum_{j} N_{j} S_{0}^{2} F_{j}(k) \frac{\sin(2kr_{j} + \phi_{j}(k))}{kr_{j}^{2}} e^{-2\sigma_{j}^{2}k^{2}} e^{-2r_{j}/\lambda_{j}(k)}$$
  
$$r_{j} = (1 + \alpha)r_{j}$$

 $r_{j} = (1 + \alpha) r_{0,j}$ •Lytle, F. W., 1999, J. Synchrotron Radiat. 6, 123 •Stumm von Bordwehr, R., 1989, Ann. Phys. (Paris) 14, 377. •EXAFS Scattering Theory: Sayers, Stern, Lytle, Phys. Rev. Lett. 27 (1971) 1204 •Rehr, Albers, Rev. Mod. Phys. 72 (2000) 621-654 •Newville, M. (2004). Fundamentals of XAFS. CARS, University of Chicago, Chicago IL, pp. 23-24 •Bunker, G. (2010). In Introduction to XAFS: A practical guide to X-ray Absorption Fine Structure Spectroscopy, pp. 92–95. CUP •Chantler et al., J Synch. Rad.19 (2012) 145-158; J Synch Rad 19 (2012) 851-862. International Tables for Crystallography Volume I

## ANALYSIS OF EXPERIMENTAL SPECTRA

- •Ab Initio Theory codes e.g. FEFF (FMS), FDMNES (FDM), GNXAS (FMS), CONTINUUM (FMS), WIEN2k (DFT)
- •Analysis and Fitting Codes e.g. IFEFFIT, MXAN, EXCURV\*, XFIT

 Different approaches useful for different materials, energy ranges. Most commonly, FDM/DFT for near-edge (XANES), MS for high energies (EXAFS)

# Different Approaches to Analysis FEFF and related codes: Spherically symmetric potentials (Muffin-tin (MT))

- FEFF and related codes: Spherically symmetric potentials (Muffin-tin (MT) approximation). Simulation of XANES spectra and fitting of EXAFS region. (developed in US) IFEFFIT (US, Australia) Athena, Artemis etc. (US)
- EXCURVE : Potentials and corresponding phase shifts are calculated for each constituent atom of the examined material sample from a superposition of neutral atomic solutions or potential files produced by *ab initio* codes are used. Simulation of XANES spectra and fitting of EXAFS region. (developed in UK)
- MXAN: Full multiple scattering calculations with MT approximation and FPMS (Non MT approximation). Structural and electronic properties are fitted in XANES region (developed in Italy)
- FDMNES: Finite Difference Method (FDM) to solve the Schrödinger equation
- (non MT approximation) and MT approximation. Pre-edge and XANES spectra are simulated (developed in France)
- Fit-it: Use FEFF and FDMNES to fit XANES (developed in Russian Federation)
- FDMX: Finite Difference Method (FDM) with full XAFS analysis. Pre-edge, XANES and XAFS spectra fitted, Inelastic mean free paths simulated and fitted, thermal broadening fitted or measured (developed in Australia from France)
- Others Shirley et al.

6. Packages Data Collection, Reduction, Prediction, Analysis 130pp? Part 6A Theoretical Prediction

**Muffin-Tin Multiple scattering** 

- 6.1 EXCURVE
- 6.2 FEFF
- 6.3 GNXAS I: PHAGEN/GNPEAK/GNXAS

Non-Muffin-Tin Multiple scattering

6.4 FPMS - Full Potential Multiple Scattering

#### Non-Muffin-Tin Finite Difference Method

- 6.5 FDMNES Finite Difference Method for Near Edge Structure
- 6.6 FDMX Finite Difference Method for XAS

#### Multiplets

6.7 TT-Multiplets

#### **Other Density Functional Theory (inc. 6.4-6.6)**

6.9 QuantumEspresso/Xspectra (First Principles Calculations with DFT)

6.10 Wien2K

- 6.11 OCEAN (Abinit, DFT) Correlation Packages
- 6.12 Exciting (DFT full potential)

#### Part 6B Data Analysis

- 6.13 Athena and Artemis
- 6.14 EDA
- 6.15 ESTRA
- 6.16 FitIt
- 6.17 GNXAS II: FITHEO
- 6.18 IFEFFIT and Larch
- 6.19 IFEFFIT with errors
- 6.20 MAX
- 6.21 MXAN
- 6.22 PRESTOPRONTO
- 6.23 RSXAP and RSFIT
- 6.24 SixPACK
- 6.25 VIPER and XANDA
- 6.26 XAFSX
- 6.27 XFIT
- 6.28 LASE